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Combinations of single-top-quark production cross-section measurements and $|f_{LV}V_{tb}|$ determinations at $\sqrt{s} = 7$ and 8 TeV with the ATLAS and CMS experiments



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ABSTRACT: This paper presents the combinations of single-top-quark production cross-section measurements by the ATLAS and CMS Collaborations, using data from LHC proton–proton collisions at $\sqrt{s} = 7$ and 8 TeV corresponding to integrated luminosities of 1.17 to 5.1 fb^{−1} at $\sqrt{s} = 7$ TeV and 12.2 to 20.3 fb^{−1} at $\sqrt{s} = 8$ TeV. These combinations are performed per centre-of-mass energy and for each production mode: t -channel, tW , and s -channel. The combined t -channel cross-sections are 67.5 ± 5.7 pb and 87.7 ± 5.8 pb at $\sqrt{s} = 7$ and 8 TeV respectively. The combined tW cross-sections are 16.3 ± 4.1 pb and 23.1 ± 3.6 pb at $\sqrt{s} = 7$ and 8 TeV respectively. For the s -channel cross-section, the combination yields 4.9 ± 1.4 pb at $\sqrt{s} = 8$ TeV. The square of the magnitude of the CKM matrix element V_{tb} multiplied by a form factor f_{LV} is determined for each production mode and centre-of-mass energy, using the ratio of the measured cross-section to its theoretical prediction. It is assumed that the top-quark-related CKM matrix elements obey the relation $|V_{td}|, |V_{ts}| \ll |V_{tb}|$. All the $|f_{LV}V_{tb}|^2$ determinations, extracted from individual ratios at $\sqrt{s} = 7$ and 8 TeV, are combined, resulting in $|f_{LV}V_{tb}| = 1.02 \pm 0.04$ (meas.) ± 0.02 (theo.). All combined measurements are consistent with their corresponding Standard Model predictions.

KEYWORDS: Hadron-Hadron scattering (experiments)

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1 Introduction

Measurements of single-top-quark production via the electroweak interaction, a process first observed in proton–antiproton ($p\bar{p}$) collisions at the Tevatron [1, 2], have entered the precision era at the Large Hadron Collider (LHC). It has become possible to measure top-quark properties using single-top-quark events [3]. Single-top-quark production is sensitive to new physics mechanisms [4] that either modify the tWb coupling [5–10] or introduce new

particles and interactions [11–16]. The production rate of single top quarks is proportional to the square of the left-handed coupling at the tWb production vertex, assuming that there are no significant tWd or tWs contributions. In the Standard Model (SM), this coupling is given by the Cabibbo-Kobayashi-Maskawa (CKM) [17, 18] matrix element V_{tb} . Indirect measurements of $|V_{tb}|$, from precision measurements of B -meson decays [19] and from top-quark decays [20–23], rely on the SM assumptions that the CKM matrix is unitary and that there are three quark generations. The most stringent indirect determination comes from a global fit to all available B -physics measurements, resulting in $|V_{tb}| = 0.999105 \pm 0.000032$ [19]. This fit also assumes the absence of any new physics mechanisms that might affect b -quarks. The most precise indirect measurement using top-quark events was performed by the CMS Collaboration in proton–proton (pp) collisions at a centre-of-mass energy of $\sqrt{s} = 7$ TeV, resulting in $|V_{tb}| = 1.007 \pm 0.016$ [23].

A direct estimate of the coupling at the tWb production vertex, $|f_{LV}V_{tb}|$, is obtained from the measured single-top-quark cross-section $\sigma_{\text{meas.}}$ and its corresponding theoretical expectation $\sigma_{\text{theo.}}$,

$$|f_{LV}V_{tb}| = \sqrt{\frac{\sigma_{\text{meas.}}}{\sigma_{\text{theo.}} (V_{tb} = 1)}}. \quad (1.1)$$

The f_{LV} term is a form factor, assumed to be real, that parameterises the possible presence of anomalous left-handed vector couplings [24]. By construction, this form factor is exactly one in the SM, while it can be different from one in models of new physics processes. The direct estimation assumes that $|V_{td}|, |V_{ts}| \ll |V_{tb}|$ [25, 26], and that the tWb interaction involves a left-handed weak coupling, like that in the SM. The $|f_{LV}V_{tb}|$ determination via single-top-quark production is independent of assumptions about the number of quark generations and the unitarity of the CKM matrix [4, 27–29]. Since the indirect determination of $|V_{tb}|$ gives a value close to unity, V_{tb} is considered equal to one in theoretical calculations of the single-top-quark cross-section. The combination of single-top-quark measurements from the Tevatron gives $|f_{LV}V_{tb}| = 1.02^{+0.06}_{-0.05}$ [30].

Single-top-quark production at a hadron collider mostly proceeds, according to the SM prediction, via three modes that can be defined at leading order (LO) in perturbative quantum chromodynamics (QCD): the exchange of a virtual W boson in the t -channel or in the s -channel, and the associated production of a top quark and a W boson (tW). Representative Feynman diagrams for these processes at LO are shown in figure 1.

In pp collisions at the LHC, the process with the largest single-top-quark production cross-section is the t -channel, where a light-flavour quark q from one of the colliding protons interacts with a b -quark by exchanging a space-like virtual W boson, producing a top quark (t -quark) and a recoiling light-flavour quark q' , called the spectator quark. For t -channel production at LO, the b -quark can be considered as directly emitted from the other proton (five-flavour-number scheme or 5FS) or it can come from gluon splitting (four-flavour-number scheme or 4FS) [31]. The kinematic properties of the spectator quark provide distinctive features for this process [32, 33]. The associated production of a W boson and a top quark has the second-largest production cross-section. In a representative process of tW production, a gluon interacts with an initial b -quark by exchanging a virtual b -quark,

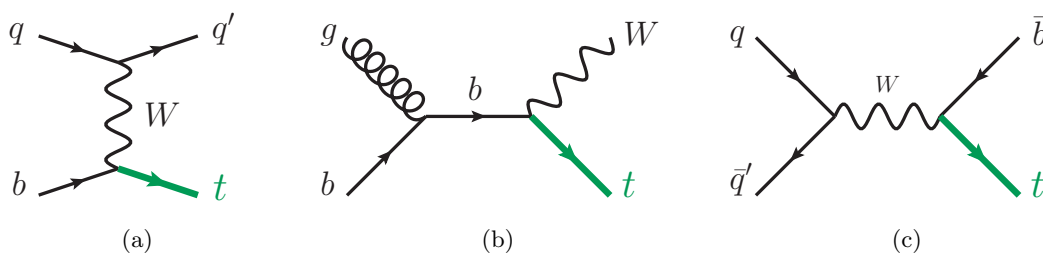


Figure 1. Representative Feynman diagrams at LO in QCD and in the five-flavour-number scheme for single-top-quark production in (a) the t -channel, (b) tW production, and (c) the s -channel.

producing a t -quark and a W boson. The measurement of this process suffers from a large background from top-quark pair ($t\bar{t}$) production [34, 35]. The s -channel cross-section is the smallest at the LHC. In this process, a quark–antiquark pair annihilates to produce a time-like virtual W boson, which decays to a t -quark and a \bar{b} -quark. This process was observed in $p\bar{p}$ collisions at the Tevatron [36] and evidence of it was reported by the ATLAS Collaboration in pp collisions at $\sqrt{s} = 8$ TeV [37].

In this paper, the t -channel, tW , and s -channel single-top-quark cross-section measurements by the ATLAS and CMS experiments are combined for each production mode, separately at pp centre-of-mass energies of 7 and 8 TeV. A combined determination of $|f_{LV}V_{tb}|$ is also presented, using as inputs the values of $|f_{LV}V_{tb}|^2$ calculated from the measured and predicted single-top-quark cross-sections in the three production modes at $\sqrt{s} = 7$ and 8 TeV. Using the same approach, results are also shown for $|f_{LV}V_{tb}|$ combinations for each production mode.

The theoretical cross-section calculations are described in section 2. Section 3 presents the cross-section measurements. The combination methodology is briefly described in section 4. Section 5 is devoted to a discussion of systematic uncertainties in the cross-section measurements as well as theoretical calculations, where the latter affect the $|f_{LV}V_{tb}|$ extraction in particular. The assumptions made about the correlation of uncertainties between the two experiments, as well as between theoretical calculations, are also discussed. Section 6 presents the combination of cross-sections for each production mode at the same centre-of-mass energy. In section 7, determinations of $|f_{LV}V_{tb}|$ are performed using all single-top-quark cross-section measurements together or by production mode. Stability tests are also shown and discussed. In section 8, the results are summarised.

2 Theoretical cross-section calculations

The theoretical predictions for the single-top-quark production cross-sections are calculated at next-to-leading order (NLO) in the strong coupling constant α_s , at NLO with next-to-next-to-leading-logarithm (NNLL) resummation (named NLO+NNLL), and at next-to-next-to-leading order (NNLO). The difference between 4FS and 5FS is small [38, 39], and the calculations use the 5FS. The NLO prediction is used in the V_{tb} combination for the t -channel and s -channel, while the NLO+NNLL prediction is used for tW , as explained

below. The NLO prediction is calculated with HATHOR (v2.1) [40, 41]. Uncertainties comprise the scale uncertainty, the α_s uncertainty, and the parton distribution function (PDF) uncertainty. The scale uncertainty is evaluated by varying the renormalisation and factorisation scales up and down together by a factor of two. The combination of the PDF+ α_s uncertainty is calculated according to the PDF4LHC prescription [42] from the envelope of the uncertainties at 68% confidence level (CL) in the MSTW2008 NLO, CT10 NLO [43], and NNPDF2.3 [44] PDF sets.

The NLO+NNLL predictions [45] are available for all single-top-quark production modes [46–48]. Uncertainties in these calculations are estimated by varying the renormalisation and factorisation scales between $m_t/2$ and $2m_t$, where m_t is the top-quark mass, and from the 90% CL uncertainties in the MSTW2008 NNLO [49, 50] PDF set. The evaluation of the PDF uncertainties is provided by the author of refs. [46–48] and is not fully compatible with the PDF4LHC prescription. The t -channel cross-sections at $\sqrt{s} = 7$ and 8 TeV are also computed at NNLO in α_s [51], with the renormalisation and factorisation scales set to m_t . This results in cross-sections which are about 0.3% and 0.6% lower than the NLO values at $\sqrt{s} = 7$ and 8 TeV respectively. However, only a limited number of scale variations are evaluated [51].

A summary of all the available theoretical cross-section predictions for t -channel, tW , and s -channel production, $\sigma_{\text{theo.}}^{t\text{-chan.}}$, $\sigma_{\text{theo.}}^{tW}$, and $\sigma_{\text{theo.}}^{s\text{-chan.}}$ respectively, with their uncertainties is shown in table 1.

In this paper, NLO predictions serve as the reference for the t - and s -channel processes, following the prescriptions presented above, because higher-order calculations and their uncertainties are not fully computed and available for the parameter values of choice. The advantage of the NLO cross-section calculations is that the configurable parameters in HATHOR can be set according to those used to generate the ATLAS and CMS simulation samples. The t - and s -channel processes do not interfere at NLO [52]. For these two processes, the entire phase space is included in the integration in order to obtain the total cross-section. The tW cross-section prediction, $\sigma_{\text{theo.}}^{tW}$, is available at NLO [41] and NLO+NNLL [47, 53]. The tW process at NLO interferes with the $t\bar{t}$ process at LO with the subsequent decay $\bar{t} \rightarrow W\bar{b}$. In the NLO prediction for tW production provided in ref. [41], a kinematic cut-off is imposed on the transverse momentum (p_T) of the outgoing b -quark, suppressing the contribution from $t\bar{t}$ production. Since the treatment of this interference in HATHOR is still being developed [54, 55], the NLO+NNLL calculation is used as reference for tW production. For the reference cross-section predictions, uncertainties corresponding to the dependence on m_t and on the LHC beam energy, E_{beam} , are evaluated. The m_t dependence is estimated by varying its central value of 172.5 GeV (the value used in the simulation samples used to measure the single-top-quark cross-sections) by ± 1 GeV, using the functional form proposed in ref. [56]. The theoretical calculations are performed at a given centre-of-mass energy while the energy of the LHC beam is measured with an uncertainty. The single-top-quark cross-sections are assumed to depend on E_{beam} according to the model given in ref. [57], with a relative uncertainty $\delta E_{\text{beam}}/E_{\text{beam}}$ of 0.1% [58]. The theoretical cross-sections that are used as reference are marked with a † in table 1.

\sqrt{s}	Process	Accuracy	$\sigma_{\text{theo.}} [\text{pb}]$
7 TeV	t -channel	NLO [†]	$63.9^{+1.9}_{-1.3} (\text{scale}) \pm 2.2 (\text{PDF}+\alpha_s) \pm 0.7 (m_t) \pm 0.1 (E_{\text{beam}})$
		NLO+NNLL	$64.6^{+2.6}_{-1.7} (\text{scale}+\text{PDF}+\alpha_s)$
		NNLO	$63.7^{+0.5}_{-0.3} (\text{scale})$
	tW	NLO	$13.2^{+0.5}_{-0.6} (\text{scale}) \pm 1.3 (\text{PDF}+\alpha_s)$
		NLO+NNLL [†]	$15.74 \pm 0.40 (\text{scale})^{+1.10}_{-1.14} (\text{PDF}+\alpha_s) \pm 0.28 (m_t) \pm 0.04 (E_{\text{beam}})$
	s -channel	NLO [†]	$4.29^{+0.12}_{-0.10} (\text{scale}) \pm 0.14 (\text{PDF}+\alpha_s) \pm 0.10 (m_t) \pm 0.01 (E_{\text{beam}})$
		NLO+NNLL	$4.63^{+0.20}_{-0.18} (\text{scale}+\text{PDF}+\alpha_s)$
8 TeV	t -channel	NLO [†]	$84.7^{+2.6}_{-1.7} (\text{scale}) \pm 2.8 (\text{PDF}+\alpha_s) \pm 0.8 (m_t) \pm 0.2 (E_{\text{beam}})$
		NLO+NNLL	$87.8^{+3.4}_{-1.9} (\text{scale}+\text{PDF}+\alpha_s)$
		NNLO	$84.2^{+0.3}_{-0.2} (\text{scale})$
	tW	NLO	$18.77^{+0.77}_{-0.82} (\text{scale}) \pm 1.70 (\text{PDF}+\alpha_s)$
		NLO+NNLL [†]	$22.37 \pm 0.60 (\text{scale}) \pm 1.40 (\text{PDF}+\alpha_s) \pm 0.38 (m_t) \pm 0.06 (E_{\text{beam}})$
	s -channel	NLO [†]	$5.24^{+0.15}_{-0.12} (\text{scale}) \pm 0.16 (\text{PDF}+\alpha_s) \pm 0.12 (m_t) \pm 0.01 (E_{\text{beam}})$
		NLO+NNLL	$5.61 \pm 0.22 (\text{scale}+\text{PDF}+\alpha_s)$

Table 1. Predicted cross-sections for single-top-quark production at $\sqrt{s} = 7$ and 8 TeV at the LHC. Uncertainties include scale and $\text{PDF}+\alpha_s$ variations, except for the NNLO predictions, which only contain the scale variation. The $\text{PDF}+\alpha_s$ uncertainties are evaluated according to the PDF4LHC prescription only for the NLO predictions. The uncertainties associated with the top-quark mass m_t and beam energy E_{beam} are also given for the NLO predictions for the t - and s -channels, and for the NLO+NNLL prediction for tW production. The value of m_t is set to 172.5 GeV in all predictions. The cross-sections marked with [†] are those used in the $|f_{\text{LV}} V_{tb}|$ combination.

3 Single-top-quark cross-section measurements at $\sqrt{s} = 7$ and 8 TeV

The t -channel single-top-quark production cross-sections, $\sigma_{t\text{-chan.}}$, were measured by the ATLAS and CMS Collaborations at $\sqrt{s} = 7$ TeV [59, 60] and 8 TeV [32, 33]. Evidence of tW production was reported at $\sqrt{s} = 7$ TeV by ATLAS [61] and CMS [62], while at $\sqrt{s} = 8$ TeV its cross-section, σ_{tW} , was measured by both experiments [34, 35]. Evidence of s -channel production was reported by ATLAS, with a measured cross-section, $\sigma_{s\text{-chan.}}$, at $\sqrt{s} = 8$ TeV [37], whereas CMS set upper limits on the s -channel production cross-section at $\sqrt{s} = 7$ and 8 TeV. The observed (expected) significance of the CMS measurement at $\sqrt{s} = 8$ TeV is 2.3 (0.8) standard deviations [63].

The ATLAS and CMS analyses use similar approaches to measure the single-top-quark production cross-sections. Both experiments select events containing at least one prompt isolated lepton (electron or muon) and at least one high- p_T jet. The analyses use various multivariate analysis (MVA) techniques, such as boosted decision trees [64–66], neural networks [67], or the matrix element method (MEM) [68, 69], to separate the signal from background. To measure the cross-section, analyses perform a binned maximum-likelihood fit to data using the distribution of the corresponding MVA discriminator. Exceptions are the ATLAS s -channel and CMS t -channel measurements at $\sqrt{s} = 8$ TeV. In the ATLAS s -channel analysis, the fit is performed simultaneously to the MEM discriminant

		ATLAS		CMS	
\sqrt{s}	Process	σ [pb]	Lumi. [fb^{-1}]	σ [pb]	Lumi. [fb^{-1}]
7 TeV	t -channel	68 ± 8	4.59	67.2 ± 6.1	1.17–1.56
	tW	16.8 ± 5.7	2.05	16^{+5}_{-4}	4.9
	s -channel	—	—	7.1 ± 8.1	5.1
8 TeV	t -channel	$89.6^{+7.1}_{-6.3}$	20.2	83.6 ± 7.8	19.7
	tW	$23.0^{+3.6}_{-3.9}$	20.3	23.4 ± 5.4	12.2
	s -channel	$4.8^{+1.8}_{-1.5}$	20.3	13.4 ± 7.3	19.7

Table 2. Summary of the single-top-quark cross-section measurements published by the ATLAS and CMS Collaborations at $\sqrt{s} = 7$ and 8 TeV. Total uncertainties are shown. Small differences between the integrated luminosity values in different analyses within the same experiment and centre-of-mass energy are due to different luminosity calibrations at the time of publication.

in the signal region and the lepton-charge distribution in the W +jets control region. The CMS t -channel measurement at $\sqrt{s} = 8$ TeV is based on a simultaneous fit to the absolute pseudorapidity (η) distributions of the recoiling light-flavour jet in events with negative and with positive lepton charge. The analyses measuring different single-top-quark production modes within the same experiment and at the same centre-of-mass energy have disjoint signal regions. Both experiments simulate the single-top-quark processes using the NLO POWHEG-BOX generator [70–74] for the matrix-element (ME) calculations. ATLAS also uses the POWHEG-BOX generator to simulate top-quark-pair background events, while CMS uses the LO MADGRAPH generator [75]. The PYTHIA [76] event generator is used for modelling the parton shower (PS), hadronisation and the underlying event in both the single-top-quark and $t\bar{t}$ processes. The cross-sections are measured assuming a value of 172.5 GeV for m_t for all top-quark processes and all centre-of-mass energies. A summary of the uncertainties in each measurement is shown in table 2, with details given in appendix A.

4 Combination methodology

The ATLAS and CMS single-top-quark production cross-section measurements shown in table 2 are combined, and the combined $|f_{LV}V_{tb}|$ value determined, using the best linear unbiased estimator (BLUE) method [77–79]. The BLUE method is applied iteratively in order to reduce a possible bias arising from the dependence of systematic uncertainties on the central value of the cross-section [80]. Convergence is reached when the central value changes by less than 0.01% compared with the previous iteration. In each iteration, the BLUE method minimises the global χ^2 by adjusting the weight for each input measurement [79]. The global χ^2 is calculated taking correlations into account. The sum of weights is required to be equal to one. Negative weights are allowed; these indicate strong correlations [81]. The number of degrees of freedom is $n - 1$, where n is number of measurements in the combination. The χ^2 and n are then used to calculate a corresponding probability [79].

The systematic uncertainties are scaled with the cross-section in each iteration, i.e. they are treated as relative uncertainties. The data and simulation statistical uncertainties are not scaled [80]. The systematic uncertainties in the s -channel cross-section combination are also not scaled because the s -channel measurements have large backgrounds.

Following the same strategy as in the input measurements by the ATLAS and CMS Collaborations, the combined cross-sections are reported at $m_t = 172.5$ GeV, not including the uncertainty associated with the m_t variation. The shift in the combined cross-section due to a variation of ± 1 GeV in the top-quark mass is given where this information is available. For the determination of the combined $|f_{LV}V_{tb}|$ value, the uncertainty in the measured cross-sections due to a variation of ± 1 GeV in the mass is considered. Uncertainties in the measurements are symmetrised, before combination, by averaging the magnitude of the downward and upward variations. More details are given in sections 5 and 6.

5 Systematic uncertainties and correlation assumptions

In order to combine single-top-quark cross-section measurements and $|f_{LV}V_{tb}|$ values, the sources of uncertainty are grouped into categories. While the categorisation and evaluation of uncertainties varies somewhat between experiments and between measurements, each individual measurement considers a complete set of uncertainties. Assumptions are made about correlations between similar sources of uncertainty in different measurements, as explained in section 5.1. Uncertainties associated with theoretical predictions are taken into account in the $|f_{LV}V_{tb}|$ combination. The correlations between similar uncertainties in different theoretical predictions are discussed in section 5.2.

5.1 Systematic uncertainties in measured cross-sections

Systematic uncertainties in the ATLAS t -channel measurements at $\sqrt{s} = 7$ and 8 TeV are evaluated using pseudoexperiments, except the background normalisation uncertainties, which are constrained in the fit to data. In the ATLAS tW measurements at $\sqrt{s} = 7$ and 8 TeV and the s -channel measurement at $\sqrt{s} = 8$ TeV, systematic uncertainties are included as nuisance parameters in profile-likelihood fits. Systematic uncertainties in the CMS t -channel and tW measurements at $\sqrt{s} = 7$ and 8 TeV are included as nuisance parameters in fits to data, except the theory modelling uncertainties in signal and backgrounds, described below, which are evaluated using pseudoexperiments. All systematic uncertainties in the CMS s -channel measurements at $\sqrt{s} = 7$ and 8 TeV are obtained through pseudoexperiments, except the background normalisation uncertainties, which are constrained in the fit to data. In the analyses where systematic uncertainties are included as nuisance parameters, the total uncertainty presented in table 2 is evaluated by varying all the nuisance parameters in the fit simultaneously. To extract the impact of each source of this type of uncertainty, these analyses use approximate procedures which neglect the correlations between sources of uncertainty introduced by the fits. Throughout this paper, individual uncertainties are taken as reported by the input analyses, regardless of the method used to determine them. The total uncertainties are evaluated as the sum in quadrature of individual contributions.

Although the sources of systematic uncertainty and the procedures used to estimate their impact on the measured cross-section are partially different in the individual analyses, it is still possible to identify contributions that describe similar physical effects. These contributions are listed below; they are grouped together, and only the resulting categories are used in the combination. Categories are treated as uncorrelated among each other. For each source of uncertainty, correlations between different measurements are assumed to be positive, unless explicitly mentioned otherwise. The stability of the cross-section and $|f_{LV}V_{tb}|$ combinations is studied by varying the correlation assumptions for the dominant uncertainties, as discussed in section 7.2.

The uncertainties in each category are listed below, with the correlation assumptions across experiments given in parentheses. These correlations correspond to those used in the cross-section combinations. They are also valid for the combination of the $|f_{LV}V_{tb}|$ extractions, unless explicitly mentioned otherwise. The symbol “—” means that the uncertainty is either considered only in the ATLAS or the CMS measurement, or is not considered at all. A summary of uncertainties in the cross-section measurements together with the corresponding correlation assumptions between experiments is provided in appendix A.

Data statistical (Correlation 0): this statistical uncertainty arises from the limited size of the data sample. It is uncorrelated between ATLAS and CMS, between production modes, and between centre-of-mass energies.

Simulation statistical (Correlation 0 and — for CMS tW at $\sqrt{s} = 7$ TeV and s -channel at $\sqrt{s} = 8$ TeV): this statistical uncertainty comes from the limited size of simulated event samples. It is uncorrelated between ATLAS and CMS, between production modes, and between centre-of-mass energies. For the CMS tW analysis at $\sqrt{s} = 7$ TeV and s -channel analysis at $\sqrt{s} = 8$ TeV, this uncertainty is evaluated as part of the total statistical uncertainty, which is also considered uncorrelated, as discussed above. More details are given in appendices A.2 and A.3.

Integrated luminosity (Correlation 0.3): this uncertainty originates from the systematic uncertainty in the integrated luminosity, as determined by the individual experiments using the methods described in refs. [82–85]. It affects the determination of both the signal and background yields. The integrated-luminosity uncertainty has a component that is correlated between ATLAS and CMS, arising from imperfect knowledge of the beam currents during van der Meer scans in the LHC accelerator [86], and an uncorrelated component from the long-term luminosity monitoring that is experiment-specific. At $\sqrt{s} = 7$ TeV, these components are 0.5% and 1.7% respectively for ATLAS and 0.5% and 2.1% respectively for CMS. At $\sqrt{s} = 8$ TeV, they are 0.6% and 1.8% respectively for ATLAS and 0.7% and 2.5% respectively for CMS. At both centre-of-mass energies, the correlation coefficient between the integrated-luminosity uncertainty in ATLAS and CMS at the same centre-of-mass energy is $\rho = 0.3$. Within the same experiment, the integrated-luminosity uncertainty is assumed to be correlated between production modes and uncorrelated between centre-of-mass energies. In section 7.2, it is shown that the combined $|f_{LV}V_{tb}|^2$ result does not depend significantly on the correlation assumptions.

Theory modelling: this category contains the uncertainties in the modelling of the simulated single-top-quark processes, as well as smaller contributions from the modelling of the $t\bar{t}$ and W +jets background processes. Both signal and background modelling are included because the uncertainties in all top-quark processes are closely related. These include initial- and final-state radiation (ISR/FSR), renormalisation and factorisation scales, NLO matching method, PS and hadronisation modelling, and PDF uncertainties. For the tW process, the uncertainty due to the treatment of interference between tW and $t\bar{t}$ final states is also included, as discussed below. These modelling uncertainties in signal and background processes are summed in quadrature in each input measurement.

- *Scales and radiation modelling* (Correlation 1)

The renormalisation and factorisation scales and ISR/FSR uncertainties account for missing higher-order corrections in the perturbative expansion and the amount of initial- and final-state radiation in simulated signal and background processes. In the ATLAS measurements of all three production modes, these uncertainties are estimated using dedicated single-top-quark and $t\bar{t}$ simulated event samples, by consistently varying the renormalisation and factorisation scales and the amount of ISR/FSR in accordance with a measurement of additional jet activity in $t\bar{t}$ events at $\sqrt{s} = 7$ TeV [87, 88]. In the ATLAS t -channel measurements, they are also estimated in W +jets simulated event samples, by varying the scale and matching parameters in the ALPGEN LO multileg generator [89] at $\sqrt{s} = 7$ TeV and by varying the parameters controlling the scale in the SHERPA LO multileg generator [90] at $\sqrt{s} = 8$ TeV. In the CMS measurements, these uncertainties are estimated by varying the renormalisation and factorisation scales, and ISR/FSR, consistently in the simulated event samples. In the CMS t -channel measurement at $\sqrt{s} = 8$ TeV, this uncertainty applies only to the signal modelling since the modelling of the dominant $t\bar{t}$ and W +jets background processes is obtained from data. However, for the t -channel analysis at $\sqrt{s} = 7$ TeV, the scales are varied in the simulated signal, $t\bar{t}$, W +jets and other single-top-quark processes. The same approach is followed in the CMS s -channel measurements at both centre-of-mass energies. The tW cross-section measurements of CMS account for this uncertainty only in the tW signal and $t\bar{t}$ background, given the negligible contributions from the W +jets and other single-top-quark processes in the dilepton final state.

Although the methods are apparently different, they mostly address the same uncertainty, hence this uncertainty is considered correlated between ATLAS and CMS. It is also considered correlated between production modes and centre-of-mass energies. The combined $|f_{LV}V_{tb}|$ result does not depend significantly on this correlation assumption, as discussed in section 7.2.

- *NLO matching* (Correlation 1 for t -channel and — for tW and s -channel)

The ATLAS measurements include an uncertainty to account for different NLO matching methods implemented in different NLO event generators. This is evaluated in single-top-quark and $t\bar{t}$ simulations by comparing the POWHEG-Box,

MC@NLO [91, 92], and MADGRAPH5_aMC@NLO [93] event generators, all interfaced to HERWIG [94] (with JIMMY [95] for the underlying-event modelling). In the CMS t -channel measurement at $\sqrt{s} = 7$ TeV, the NLO matching uncertainty is evaluated by comparing POWHEG-BOX with CompHEP [96, 97]. In the CMS t -channel analysis at $\sqrt{s} = 8$ TeV, this uncertainty accounts for different NLO matching methods in the t -channel signal event generator, as well as for differences between event generation in the 4FS and 5FS, by comparing POWHEG-BOX with MADGRAPH. The NLO matching uncertainty is considered correlated between ATLAS and CMS, between production modes, and between centre-of-mass energies. In the CMS tW and s -channel analyses at $\sqrt{s} = 7$ and 8 TeV, this uncertainty is not considered, since the modelling uncertainties in the scheme to remove overlap with $t\bar{t}$ are dominant in the tW analysis and the renormalisation/factorisation scale is dominant in the s -channel analysis. The results of the stability test for this uncertainty are shown in section 7.2.

- *Parton shower and hadronisation* (Correlation 1)

In both experiments, the difference between the PYTHIA and HERWIG showering programs is considered in the jet energy scale (JES) [98–101] and b -tagging calibration [102–106]. The ATLAS analyses additionally include an uncertainty in the PS and hadronisation modelling in simulated single-top-quark and $t\bar{t}$ events, evaluated by comparing the POWHEG-BOX event generator interfaced to PYTHIA or to HERWIG. The CMS analyses additionally include an uncertainty in the $t\bar{t}$ and W +jets backgrounds estimated with the MADGRAPH event generator interfaced to PYTHIA. It is evaluated in simulated event samples where the value of the ME/PS matching threshold in the MLM method [107] is doubled or halved from its initial value. The CMS t -channel measurement at $\sqrt{s} = 8$ TeV does not consider this uncertainty in the $t\bar{t}$ and W +jets backgrounds since the distribution and normalisation of the $t\bar{t}$ and W +jets processes are derived mostly from data. In the CMS tW analyses at $\sqrt{s} = 7$ and 8 TeV, the contributions of the W +jets and other single-top-quark processes in the dilepton final state are negligible.

This uncertainty is considered correlated between ATLAS and CMS, between different production modes, and between different centre-of-mass energies. The combined $|f_{LV}V_{tb}|$ result does not depend significantly on this correlation assumption, as shown in section 7.2.

- *Parton distribution functions* (Correlation 1)

The PDF uncertainty is evaluated following the PDF4LHC procedures [42, 108, 109] and is considered correlated between ATLAS and CMS, between different production modes, and between different centre-of-mass energies.

- *tW and $t\bar{t}$ interference* (Correlation 1 for tW and — for t - and s -channels)

The tW process interferes with $t\bar{t}$ production at NLO [110–112]. In both ATLAS and CMS, two simulation approaches are compared: diagram removal (DR) [110] and diagram subtraction (DS) [27, 110]. In the DR approach, all NLO diagrams that overlap

with the doubly resonant $t\bar{t}$ contributions are removed from the calculation of the tW amplitude. This approach accounts for the interference term, but it is not gauge invariant (though the effect is numerically negligible) [110]. In the DS approach, a subtraction term is built into the amplitude to cancel out the $t\bar{t}$ component close to the top-quark resonance while respecting gauge invariance.

The DR approach is the default, and the comparison with the DS approach is used to assess this systematic uncertainty. For the tW analyses, this uncertainty is considered correlated between the two experiments and between different centre-of-mass energies.

- *Modelling of the top-quark p_T spectrum* (Correlation —)

In the CMS tW and s -channel analyses at $\sqrt{s} = 8$ TeV, the simulated $t\bar{t}$ events are reweighted to correct the p_T spectrum of the generated top quarks, which was found to be significantly harder than the spectrum observed in data in differential cross-section measurements [113, 114]. To estimate the uncertainty related to this mismodelling, the tW measurement is repeated without the reweighting, and the change relative to the default result is taken as the uncertainty. In the CMS s -channel analysis, the measurement is repeated with the effect of the weights removed and doubled. The resulting variation in the cross-section is symmetrised. This uncertainty is not considered in the CMS t -channel measurement at $\sqrt{s} = 8$ TeV where the modelling of the $t\bar{t}$ background is extracted from data. In the ATLAS measurements, modelling uncertainties in the top-quark p_T spectrum in $t\bar{t}$ events [115] are covered by the PS and hadronisation uncertainty and they are found to be small in comparison with other systematic uncertainties. This uncertainty is considered correlated between the CMS tW and s -channel analyses at $\sqrt{s} = 8$ TeV.

- *Dependence on the top-quark mass* (Correlation 1)

The measured single-top-quark cross-sections shown in table 2 assume a nominal m_t value of 172.5 GeV. The dependence of the measured cross-section on m_t is estimated for the ATLAS t -channel measurements at $\sqrt{s} = 7$ and 8 TeV and for the ATLAS tW measurement at $\sqrt{s} = 8$ TeV. It is determined using dedicated simulations of single-top-quark and $t\bar{t}$ samples with different m_t values. The cross-section measurements assuming the different m_t values are interpolated using a first- or a second-order polynomial, for which the constant term is given by the central value of $m_t = 172.5$ GeV. The CMS measurements at $\sqrt{s} = 8$ TeV provide information for a variation of ± 2 GeV in the top-quark mass, which is scaled to a ± 1 GeV shift assuming a linear dependence. For the CMS t -channel and tW measurements at $\sqrt{s} = 8$ TeV, the changes in cross-sections are symmetrised and reported as uncertainties. In the CMS s -channel analysis, the change in the cross-section is determined for the up and down variation of m_t . No estimates are available for the CMS t -channel analysis at $\sqrt{s} = 7$ TeV, the ATLAS and CMS tW analyses at $\sqrt{s} = 7$ TeV or the ATLAS s -channel analysis at $\sqrt{s} = 8$ TeV. The top-quark-mass uncertainty is small

for each measurement, thus the impact of not evaluating it for these measurements is negligible.

In this paper, a symmetrised uncertainty in the measured cross-section due to a variation of ± 1 GeV in the top-quark mass is considered. When the full cross-section dependence on the top-quark mass is available for a given production mode at a given centre-of-mass energy, the sign of the dependence of the uncertainty per unit of mass is taken into account in the correlations. In the case of the CMS t -channel and tW measurements at $\sqrt{s} = 8$ TeV, where the sign of the dependence is not available, it is assumed that the sign is the same as for the ATLAS measurement, since the phase space and background composition are comparable between CMS and ATLAS. Given that the uncertainty in the measured cross-section is considered for the same m_t variation and considering the sign of the dependence when available, this uncertainty is considered correlated between ATLAS and CMS and between different centre-of-mass energies and uncorrelated between the t -channel and tW production modes.

Background normalisation (Correlation 0): three background uncertainties are considered: in top-quark background ($t\bar{t}$ and other single-top-quark processes), in other background determined from simulation (W/Z +jets, diboson, and other smaller background channels), and in background estimated from data (multijet background from misidentified and non-prompt leptons). The exceptions are the t -channel measurements at $\sqrt{s} = 7$ TeV, where the background from simulation includes top-quark background, as shown in tables 9–13 in appendix A. The normalisation of the main background processes is determined from data, either by inclusion of normalisation uncertainties as nuisance parameters in the fit used to extract the signal, or through dedicated techniques based on data. In the t -channel and s -channel measurements, the uncertainties in the theoretical cross-section predictions for the top-quark, W/Z +jets, and diboson processes are included. In the tW measurements, the uncertainties in the theoretical cross-section predictions for the top-quark and diboson processes are taken into account. In the ATLAS measurements of the t -channel process at $\sqrt{s} = 7$ and 8 TeV, the uncertainty in the multijet background is estimated by comparing background estimates made using different techniques based on simulation and data samples. In the ATLAS tW analyses at $\sqrt{s} = 7$ and 8 TeV, the normalisation uncertainty in the background from misidentified and non-prompt leptons is obtained from variations in the data-based estimate. In the ATLAS s -channel analysis, the uncertainty assigned to the normalisation of the multijet background is based on control samples. For all CMS measurements, background normalisations are constrained in the fits to data. In the CMS measurements of the t -channel and s -channel processes, the uncertainties in the multijet background are assessed by comparing the results of alternative background estimation methods based on data. Hence, the associated uncertainties are considered uncorrelated between ATLAS and CMS, between different production modes, and between different centre-of-mass energies.

Jets: in the analyses, the uncertainties related to the reconstruction and energy calibration of jets are propagated through variations in the modelling of the detector response.

These uncertainties, classified in categories as JES, jet identification (JetID), and jet energy resolution (JER), are discussed below.

- *Jet energy scale* (Correlation 0 and — for JES flavour)

The JES is derived using information from data and simulation. Its uncertainty increases with increasing $|\eta|$ and decreases with increasing p_T of the reconstructed jet.

For all of the ATLAS measurements, except the tW measurement at $\sqrt{s} = 7$ TeV, the JES uncertainty is split into components originating from the jet calibration procedure; most of them are derived from in situ techniques based on data [98, 99]. These components are categorised as modelling, detector, calibration method, and statistical components, which are grouped into the “JES common” uncertainty, as well as a flavour-dependence component (“JES flavour”), which accounts for the flavour composition of the jets and the calorimeter response to jets of different flavours. The modelling of additional pp collisions in each bunch-crossing (pile-up) is considered separately, as discussed below. The η -dependent component is dominant for the t -channel production mode. Thus, the JES common uncertainty is considered uncorrelated between the t -channel and the other single-top-quark production modes. For the tW analysis at $\sqrt{s} = 8$ TeV, the modelling component, which is constrained in the fit to data, is dominant. The uncertainty in the flavour composition of the jets is dominant for the s -channel.

For the CMS measurements, sources contributing to the JES uncertainty are combined together into the “JES common” uncertainty, and the effect is propagated to the cross-section measurements through η - and p_T -dependent JES uncertainties [100, 101]. The jet energy corrections and their corresponding uncertainties are extracted from data. The JES uncertainty is estimated from its effect on the normalisation and shape of the discriminant in each analysis. The JES uncertainty is considered uncorrelated between the t -channel and the other single-top-quark production modes because it is dominated by the forward jet in the t -channel.

The correlation between the JES common uncertainty (or the JES uncertainty for the tW measurement at $\sqrt{s} = 7$ TeV) in ATLAS and the JES uncertainty in CMS follows the prescription in refs. [116, 117], with the slight differences for the t -channel described above. The JES common (or JES) uncertainty is considered uncorrelated between ATLAS and CMS, between centre-of-mass energies, and between production modes. Within the ATLAS experiment, the JES common uncertainty is considered correlated between tW and s -channel and uncorrelated between t -channel and the other production modes. For the ATLAS t -channel analyses, a correlation of 0.75 is assumed between $\sqrt{s} = 7$ and 8 TeV, since these analyses are mainly affected by the same uncertainty components. This correlation value is estimated by comparing variations of the JES uncertainty components in these two measurements.

In all CMS measurements and in the ATLAS tW measurement at $\sqrt{s} = 7$ TeV, the JES uncertainty is not split and therefore the JES flavour uncertainty is included in the overall JES uncertainty. For the ATLAS measurements where this component

is available, the JES flavour uncertainty is considered correlated between different production modes and uncorrelated between centre-of-mass energies.

The JES uncertainty is one of the dominant contributions in most of the single-top-quark measurements. To ensure the robustness of the results against the correlation assumptions for this large uncertainty, the combination is performed with alternative correlation values, as discussed in section 7.2.

- *Jet identification* (Correlation —)

In the ATLAS measurements, the JetID uncertainty includes the jet and vertex reconstruction efficiency uncertainties. In the CMS measurements, this uncertainty is included in the JES uncertainty. For ATLAS, it is considered correlated between the different production modes at the same centre-of-mass energy and uncorrelated for the other cases.

- *Jet energy resolution* (Correlation 0)

The uncertainty in the JER, which is not split into components, is extracted from data. Generally, the JER uncertainty is propagated via a nuisance parameter in the signal extraction fit, except for the ATLAS t -channel measurements at $\sqrt{s} = 7$ and 8 TeV, and the CMS s -channel measurement, where this uncertainty is determined using pseudoexperiments. The JER uncertainty is considered uncorrelated between ATLAS and CMS, and between centre-of-mass energies. It is considered correlated between different production modes.

Detector modelling: this category includes the uncertainty in the modelling of leptons, magnitude of the missing transverse momentum (E_T^{miss}), and identification of jets from b -quarks (b -tagging).

- *Lepton modelling* (Correlation 0)

The lepton modelling uncertainty includes components associated with the lepton energy scale and resolution, reconstruction and trigger efficiencies. This uncertainty is considered uncorrelated between ATLAS [118–121] and CMS [122] and between different centre-of-mass energies, since it is determined from data. It is considered correlated between different production modes.

- *Hadronic part of the high-level trigger* (Correlation —)

In the CMS t -channel cross-section measurement at $\sqrt{s} = 7$ TeV, the high-level trigger (HLT) criteria for the electron channel are based on the presence of an electron together with a b -tagged jet. In this analysis, the uncertainty in the modelling of the hadronic part of the HLT requirement is determined from data. This uncertainty is only evaluated in this one measurement.

- E_T^{miss} modelling (Correlation 0)

The ATLAS measurements include separate components for the uncertainties in the energy scale and resolution of the E_T^{miss} [123]. The CMS measurements account for a

combined E_T^{miss} scale and resolution uncertainty [100, 124], arising from the jet-energy uncertainties. Additionally, CMS accounts for an uncertainty in E_T^{miss} arising from energy deposits in the detector that are not included in the reconstruction of leptons, photons, and jets. The E_T^{miss} uncertainty is considered uncorrelated between ATLAS and CMS, and between different centre-of-mass energies. It is considered correlated between production modes, except for the ATLAS and CMS tW analyses at $\sqrt{s} = 8$ TeV, where it is considered uncorrelated with the other production modes because the E_T^{miss} uncertainty is constrained in the fit to data. In the ATLAS tW analysis at $\sqrt{s} = 7$ TeV, this uncertainty is included in the pile-up modelling uncertainty.

- *b-tagging* (Correlation 0)

In the ATLAS analyses, b -tagging modelling uncertainties are split into components associated with b -quark, c -quark, and light-flavour quark and gluon jets [102–104]. They are evaluated by varying the p_T -dependence (η -dependence in the case of light-flavour jets) of the flavour-dependent scale factors applied to each jet in simulation within a range that reflects the systematic uncertainty in the measured tagging efficiency and misidentification rates. This uncertainty is not considered in the ATLAS tW analysis at $\sqrt{s} = 7$ TeV because no b -tagging criterion is applied in the event selection. In the CMS measurements, the uncertainties in b -tagging efficiency and misidentification rates of jets initiated by light-flavour quarks and gluons are derived from data, using control samples [105, 106]. The CMS uncertainties are propagated to the cross-section measurements using pseudoexperiments. Exceptions are the t -channel measurement at $\sqrt{s} = 7$ TeV and the tW measurement at $\sqrt{s} = 8$ TeV, where these uncertainties are constrained in the fit to data.

The two collaborations split up the different sources of systematic uncertainties related to b -tagging in a different way. However, the different sources are combined by adding their contributions in quadrature to obtain a single b -tagging uncertainty per analysis. This means that the b -tagging uncertainty also contains the uncertainties associated with the misidentification rates of jets initiated by charm quarks, light-flavour quarks and gluons. The resulting uncertainty is considered uncorrelated between ATLAS and CMS, and between different centre-of-mass energies. It is considered correlated between different production modes.

- *Pile-up modelling* (Correlation 0)

In both ATLAS and CMS, simulated events are reweighted to match the distribution of the average number of interactions per bunch-crossing in data. The corresponding uncertainty is obtained from in situ techniques based on data and simulated event samples. In the ATLAS analyses at $\sqrt{s} = 7$ TeV, the uncertainty due to pile-up is derived from the impact of the reweighting on E_T^{miss} . In the ATLAS analyses at $\sqrt{s} = 8$ TeV, this uncertainty is evaluated as a component of the JES, separated into four terms (number of primary vertices, average number of collisions per bunch-crossing, average pile-up energy density in the calorimeter, and p_T dependence) since the pile-up calibration (assuming average conditions during 8 TeV data-taking) is

applied to both data and simulation before selecting and calibrating the jets [117]. In CMS, the reweighting uses a model with a free parameter that can be interpreted as an effective cross-section for inelastic pp interactions. This uncertainty is obtained from a fit to the number of additional primary vertices in simulation. In the CMS analyses, this uncertainty is introduced as a nuisance parameter in the fit. The only exception is the s -channel measurement, where the pile-up uncertainty is estimated from pseudoexperiments. In all cases, the effects of pile-up on the jet energy and the isolation of leptons are taken into account in the jet and lepton uncertainties respectively. The pile-up uncertainty is considered uncorrelated between ATLAS and CMS and between different centre-of-mass energies. It is considered correlated between different production modes [116, 117].

5.2 Systematic uncertainties in theoretical cross-section predictions

The systematic uncertainties in the combined $|f_{LV}V_{tb}|$ value are evaluated from uncertainties in the individual cross-section measurements $\sigma_{\text{meas.}}$ and the theoretical predictions $\sigma_{\text{theo.}}$. The uncertainties associated with $\sigma_{\text{theo.}}$ are discussed in section 2; they are summarised in table 1. The correlation assumptions for the systematic uncertainties related to the theoretical cross-section are explained below. In section 7.2, the stability of the $|f_{LV}V_{tb}|$ combination against variations in the correlations is examined. For clarity, the correlations are given in parentheses next to the systematic-uncertainty name. These correlations are used in the combination of the $|f_{LV}V_{tb}|$ extractions.

PDF+ α_s (Correlation 1 for centre-of-mass energies and 0.5 for production modes): the PDF uncertainty is considered correlated between centre-of-mass energies and 50% correlated between production modes, since different production modes have one initial-state particle in common (a quark or a gluon), but not both.

Renormalisation and factorisation scales (Correlation 1 for t -channel and s -channel and 0 for tW): the renormalisation and factorisation scale uncertainties in $\sigma_{\text{theo.}}$ are considered correlated between production modes and centre-of-mass energies, except between the tW production mode and the other production modes, where they are considered uncorrelated because the tW prediction is computed at a different order in perturbation theory.

Top-quark mass (Correlation 1): the uncertainty due to m_t is evaluated by varying m_t from its central value of 172.5 GeV by ± 1 GeV and evaluating the corresponding change in cross-section using the parameterisation given in ref. [56], as discussed in section 2. This uncertainty is considered correlated between centre-of-mass energies and production modes.

E_{beam} (Correlation 1): the uncertainty in the cross-section due to the uncertainty in E_{beam} is estimated by computing the cross-section variation corresponding to a ± 1 standard deviation shift in the beam-energy uncertainty. It is considered correlated between centre-of-mass energies and production modes.

6 Combinations of cross-section measurements

The cross-section measurements described in section 3 are combined at each centre-of-mass energy for each production mode. Systematic uncertainties are categorised and correlation assumptions are employed according to section 5. The combinations are performed using the iterative BLUE method, as described in section 4.

As discussed in section 4, the uncertainty in the measured cross-section associated with the m_t variation is not considered in the combination of cross-sections. However, the shift in the combined cross-section resulting from a variation of ± 1 GeV in the top-quark mass is provided where this information is available. This is calculated by repeating the combination with the up-shifted and down-shifted input cross-sections. In measurements where only the magnitude of the shift is available for one experiment, the sign of the shift is assumed to be the same for both experiments, as discussed in section 5.1. If the uncertainty associated with the m_t variation is not available for one or both of the input measurements, then no shift in the combined cross-section is given.

Additional information about the uncertainties considered in the combination of cross-section measurements is provided in appendix A.

6.1 Combinations of t -channel cross-section measurements

The combination of the ATLAS and CMS t -channel cross-section measurements at $\sqrt{s} = 7$ TeV [59, 60] results, after one iteration, in

$$\sigma_{t\text{-chan.}} = 67.5 \pm 2.4 \text{ (stat.)} \pm 5.0 \text{ (syst.)} \pm 1.1 \text{ (lumi.) pb} = 67.5 \pm 5.7 \text{ pb.}$$

The relative uncertainty is 8.4%, which improves on the uncertainty of 9.1% in the most precise individual measurement from CMS [60]. The χ^2 for the combination is 0.01, corresponding to a probability of 93%. The CMS weight in the combination is 0.58, while the ATLAS weight is 0.42. The overall correlation between the two measurements is 20%. The contribution from each uncertainty category to the total uncertainty in the combined t -channel cross-section measurement at $\sqrt{s} = 7$ TeV is shown in table 3(a).

The combination of the ATLAS and CMS t -channel cross-section measurements at $\sqrt{s} = 8$ TeV [32, 33] results, after two iterations, in a cross-section of

$$\sigma_{t\text{-chan.}} = 87.7 \pm 1.1 \text{ (stat.)} \pm 5.5 \text{ (syst.)} \pm 1.5 \text{ (lumi.) pb} = 87.7 \pm 5.8 \text{ pb.}$$

The relative uncertainty is 6.7%, which improves on the uncertainty of 7.5% in the most precise individual measurement from ATLAS [32]. The χ^2 for the combination is 0.59, corresponding to a probability of 44%. This probability is lower than the probability of the combination at $\sqrt{s} = 7$ TeV because of the differences between the ATLAS and CMS measured cross-sections and their small uncertainties. The ATLAS weight in the combination is 0.68, while the CMS weight is 0.32. The overall correlation between the two measurements is 42%. This is larger than the correlation between the measurements at $\sqrt{s} = 7$ TeV because the statistical and detector uncertainties are lower, thus increasing the importance of the theory modelling uncertainty (which is correlated between the two

(a)			(b)		
$\sigma_{t\text{-chan.}}, \sqrt{s} = 7 \text{ TeV}$			$\sigma_{t\text{-chan.}}, \sqrt{s} = 8 \text{ TeV}$		
Combined cross-section	67.5 pb		Combined cross-section	87.7 pb	
Uncertainty category	Uncertainty		Uncertainty category	Uncertainty	
	[%]	[pb]		[%]	[pb]
Data statistical	3.5	2.4	Data statistical	1.3	1.1
Simulation statistical	1.4	0.9	Simulation statistical	0.6	0.5
Integrated luminosity	1.7	1.1	Integrated luminosity	1.7	1.5
Theory modelling	5.1	3.5	Theory modelling	5.3	4.7
Background normalisation	1.9	1.3	Background normalisation	1.2	1.1
Jets	3.4	2.3	Jets	2.6	2.3
Detector modelling	3.4	2.3	Detector modelling	1.8	1.6
Total syst. unc. (excl. lumi.)	7.5	5.0	Total syst. unc. (excl. lumi.)	6.3	5.5
Total syst. unc. (incl. lumi.)	7.6	5.2	Total syst. unc. (incl. lumi.)	6.5	5.7
Total uncertainty	8.4	5.7	Total uncertainty	6.7	5.8

Table 3. Contribution from each uncertainty category to the combined t -channel cross-section ($\sigma_{t\text{-chan.}}$) uncertainty at (a) $\sqrt{s} = 7 \text{ TeV}$ and (b) $\sqrt{s} = 8 \text{ TeV}$. The total uncertainty is computed by adding in quadrature all the individual systematic uncertainties (including the uncertainty in the integrated luminosity) and the statistical uncertainty in data. Correlations of systematic uncertainties between experiments are presented in appendix A.1.

experiments), as shown in appendix A.1. The contribution from each uncertainty category to the total uncertainty in the combined t -channel cross-section measurement at $\sqrt{s} = 8 \text{ TeV}$ is shown in table 3(b).

At both centre-of-mass energies, the uncertainties from theory modelling are found to be dominant. Details of the central values, the impact of individual sources of uncertainties, and their correlations between experiments at $\sqrt{s} = 7$ and 8 TeV can be found in appendix A.1.

The shift in the combined cross-section at $\sqrt{s} = 8 \text{ TeV}$ from a variation of $\pm 1 \text{ GeV}$ in the top-quark mass is $\mp 0.8 \text{ pb}$, which is similar to the shifts in the input measurements for the same m_t variation. The shift in the combined cross-section at $\sqrt{s} = 7 \text{ TeV}$ is not evaluated since no estimate is available for the CMS input measurement at $\sqrt{s} = 7 \text{ TeV}$.

6.2 Combinations of tW cross-section measurements

The combination of the ATLAS and CMS tW cross-section measurements at $\sqrt{s} = 7 \text{ TeV}$ [61, 62] yields, after two iterations, a cross-section of

$$\sigma_{tW} = 16.3 \pm 2.3 \text{ (stat.)} \pm 3.3 \text{ (syst.)} \pm 0.7 \text{ (lumi.) pb} = 16.3 \pm 4.1 \text{ pb.}$$

The relative uncertainty is 25%, which improves on the uncertainty of 28% in the most precise individual measurement from CMS [62]. The χ^2 for the combination is 0.01, corresponding to a probability of 91%. The CMS weight in the combination is 0.59, while the

(a)			(b)		
$\sigma_{tW}, \sqrt{s} = 7 \text{ TeV}$			$\sigma_{tW}, \sqrt{s} = 8 \text{ TeV}$		
Combined cross-section	16.3 pb		Combined cross-section	23.1 pb	
Uncertainty category	Uncertainty		Uncertainty category	Uncertainty	
	[%]	[pb]		[%]	[pb]
Data statistical	14.0	2.3	Data statistical	4.7	1.1
Simulation statistical	0.8	0.1	Simulation statistical	0.8	0.2
Integrated luminosity	4.4	0.7	Integrated luminosity	3.6	0.8
Theory modelling	13.9	2.3	Theory modelling	11.8	2.7
Background normalisation	6.0	1.0	Background normalisation	2.2	0.5
Jets	11.5	1.9	Jets	6.2	1.4
Detector modelling	6.2	1.0	Detector modelling	4.9	1.1
Total syst. unc. (excl. lumi.)	20.0	3.3	Total syst. unc. (excl. lumi.)	14.4	3.3
Total syst. unc. (incl. lumi.)	20.5	3.3	Total syst. unc. (incl. lumi.)	14.8	3.4
Total uncertainty	24.8	4.1	Total uncertainty	15.6	3.6

Table 4. Contribution from each uncertainty category to the combined tW cross-section (σ_{tW}) uncertainty at (a) $\sqrt{s} = 7 \text{ TeV}$ and (b) $\sqrt{s} = 8 \text{ TeV}$. The total uncertainty is computed by adding in quadrature all the individual systematic uncertainties (including the uncertainty in the integrated luminosity) and the statistical uncertainty in data. Correlations of systematic uncertainties between experiments are presented in appendix A.2.

ATLAS weight is 0.41. The overall correlation between the two measurements is 17%. The contribution from each uncertainty category to the total uncertainty in the combined tW cross-section measurement at $\sqrt{s} = 7 \text{ TeV}$ is shown in table 4(a).

The combination of the ATLAS and CMS tW cross-section measurements at $\sqrt{s} = 8 \text{ TeV}$ [34, 35] results, after two iterations, in

$$\sigma_{tW} = 23.1 \pm 1.1 \text{ (stat.)} \pm 3.3 \text{ (syst.)} \pm 0.8 \text{ (lumi.) pb} = 23.1 \pm 3.6 \text{ pb.}$$

The relative uncertainty is 15.6%, which improves on the uncertainty of 16.5% in the most precise individual measurement from ATLAS [34]. The χ^2 for the combination is 0.01, corresponding to a probability of 94%. The ATLAS weight in the combination is 0.70, while the CMS weight is 0.30. The overall correlation between the two measurements is 40%. Similar to the t -channel, this is larger than the correlation between the measurements at $\sqrt{s} = 7 \text{ TeV}$ due to the increased importance of the theory modelling uncertainties. The contribution from each uncertainty category to the total uncertainty in the combined tW cross-section measurement at $\sqrt{s} = 8 \text{ TeV}$ is shown in table 4(b).

At both centre-of-mass energies, the uncertainties in the theory modelling are found to be dominant. The jet uncertainties are also important. Details of the central values, the impact of individual sources of uncertainties, and their correlations between experiments at $\sqrt{s} = 7$ and 8 TeV are presented in appendix A.2.

$\sigma_{s\text{-chan.}}, \sqrt{s} = 8 \text{ TeV}$		
Combined cross-section	4.9 pb	
Uncertainty category	Uncertainty	
	[%]	[pb]
Data statistical	16	0.8
Simulation statistical	12	0.6
Integrated luminosity	5	0.2
Theory modelling	14	0.7
Background normalisation	8	0.4
Jets	13	0.6
Detector modelling	8	0.4
Total syst. unc. (excl. lumi.)	25	1.2
Total syst. unc. (incl. lumi.)	25	1.2
Total uncertainty	30	1.4

Table 5. Contribution from each uncertainty category to the combined s -channel cross-section ($\sigma_{s\text{-chan.}}$) uncertainty at $\sqrt{s} = 8 \text{ TeV}$. The total uncertainty is computed by adding in quadrature all the individual systematic uncertainties (including the uncertainty in the integrated luminosity) and the statistical uncertainty in data. Correlations of systematic uncertainties between experiments are presented in appendix A.3.

The shift in the combined cross-section at $\sqrt{s} = 8 \text{ TeV}$ from a variation of $\pm 1 \text{ GeV}$ in the top-quark mass is $\pm 1.1 \text{ pb}$, which is similar in magnitude to that in the input measurements for the same m_t variation. The shift in the combined cross-section at $\sqrt{s} = 7 \text{ TeV}$ is not evaluated since no estimates are available for the input measurements at $\sqrt{s} = 7 \text{ TeV}$.

6.3 Combination of s -channel cross-section measurements

The ATLAS and CMS s -channel cross-section measurements suffer from large backgrounds, and the cross-section measurements have large uncertainties. Since the systematic uncertainties mainly affect the background prediction, they are not scaled in the iterative BLUE procedure. Only the luminosity uncertainty is scaled with the central value. The combination of the ATLAS and CMS s -channel cross-section measurements at $\sqrt{s} = 8 \text{ TeV}$ [37, 63] results, after two iterations, in a cross-section of

$$\sigma_{s\text{-chan.}} = 4.9 \pm 0.8 \text{ (stat.)} \pm 1.2 \text{ (syst.)} \pm 0.2 \text{ (lumi.) pb} = 4.9 \pm 1.4 \text{ pb.}$$

The relative uncertainty is 30%, very similar to the most precise individual measurement from ATLAS [37]. The χ^2 for the combination is 1.45, corresponding to a probability of 23%. The ATLAS weight in the combination is 0.99, while the CMS weight is 0.01. The overall correlation between the two measurements is 15%. The contribution from each uncertainty category to the total uncertainty in the combined s -channel cross-section measurement at $\sqrt{s} = 8 \text{ TeV}$ is shown in table 5.

Since the ATLAS measurement has a large weight in the combination, the importance of each uncertainty in the combination is similar to that in the ATLAS measurement, as presented in appendix A.3.

The shift in the combined cross-section at $\sqrt{s} = 8$ TeV from a variation in the top-quark mass is not evaluated since no estimate is available for the ATLAS input measurement.

6.4 Summary of cross-section combinations

A summary of the cross-sections measured by ATLAS and CMS and their combinations in all single-top-quark production modes at each centre-of-mass energy is shown in figure 2. The measurements are compared with the theoretical predictions shown in table 1: NNLO for t -channel only, NLO and NLO+NNLL for all three production modes. For the NLO calculation, the renormalisation- and factorisation-scale uncertainties and the sum in quadrature of the contributions from scale, PDF, and α_s are shown separately. Only the scale uncertainty is shown for the NNLO calculation. For the NLO+NNLL calculation, the sum in quadrature of the contributions from scale, PDF, and α_s is shown. All measurements are in good agreement with their corresponding theoretical predictions within their total uncertainties.

The stability of the combinations of the cross-section measurements to variations in the correlation assumptions, discussed in section 5, is checked for the theory modelling, JES, the most important contributions to the theoretical cross-section predictions (i.e. PDF+ α_s and scale) and the integrated luminosity. The results of these tests show that their impacts on the cross-section combinations are very small, similar to the stability tests for the combination of the $|f_{LV}V_{tb}|$ values discussed in section 7.2.

7 Combinations of $|f_{LV}V_{tb}|$ determinations

The measured cross-section for a given single-top-quark production mode, $\sigma_{\text{meas.}}$, has a linear dependence on $|f_{LV}V_{tb}|^2$ as defined in eq. (1.1). Thus, a value of $|f_{LV}V_{tb}|^2$ is extracted from each cross-section measurement and the corresponding theoretical prediction (presented in sections 3 and 2 respectively). These values are then combined per channel, and in an overall $|f_{LV}V_{tb}|^2$ combination. In the overall combination, the value from the CMS measurement of $\sigma_{s\text{-chan.}}$ is excluded. The reason for excluding the CMS s -channel analysis from the overall $|f_{LV}V_{tb}|^2$ combination is that, at the same centre-of-mass energy, the CMS t -channel determination has strong correlations with the s -channel determination, which contains relatively large uncertainties. The strong correlation between these two measurements makes the combined $|f_{LV}V_{tb}|^2$ value strongly dependent on the correlation assumptions for the dominant uncertainties. This results in a large variation of the combined $|f_{LV}V_{tb}|^2$ value for different correlation assumptions.

All uncertainties in $\sigma_{\text{meas.}}$ and $\sigma_{\text{theo.}}$ are propagated to the $|f_{LV}V_{tb}|^2$ values, taking into account the correlations described in section 5. The combined value of $|f_{LV}V_{tb}|^2$ is evaluated using the reference theoretical cross-section central values marked with a † in table 1, where it can also be seen that the E_{beam} uncertainty is negligible compared to other uncertainties. For the most precise measurements (i.e. for $\sigma_{t\text{-chan.}}$ cross-section

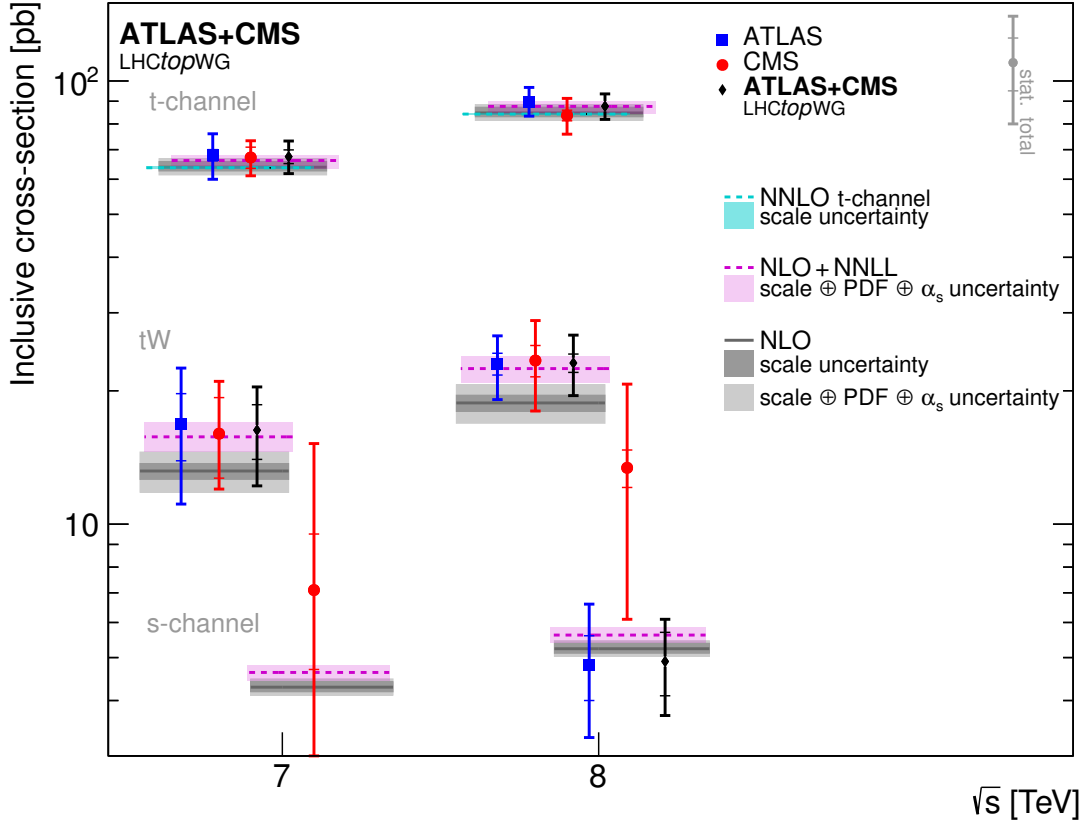


Figure 2. Single-top-quark cross-section measurements performed by ATLAS and CMS, together with the combined results shown in sections 6.1–6.3. These measurements are compared with the theoretical predictions at NLO and NLO+NNLL for all three production modes and the prediction at NNLO for t -channel only. The corresponding theoretical uncertainties are also presented. The scale uncertainty for the NNLO prediction is small and is presented as a narrow band under the dashed line.

measurements at $\sqrt{s} = 8$ TeV), which have a large expected impact on the combination, the other theoretical calculations from table 1 are used as cross-checks.

Table 6 contains a summary of the individual $|f_{LV}V_{tb}|^2$ determinations that are the inputs to the overall $|f_{LV}V_{tb}|^2$ combination, together with their experimental and theoretical uncertainties using the reference theoretical cross-sections and uncertainties. For the same processes and at the same centre-of-mass energies, there are some important differences between uncertainty categories. In analyses based on t -channel events at $\sqrt{s} = 7$ TeV, the data statistical uncertainty is larger in CMS than in ATLAS because the two experiments use data samples of different size. Differences in the category of jet uncertainties are due to the evaluation of the JES uncertainty in ATLAS using pseudoexperiments, while this uncertainty is introduced as a nuisance parameter in the fit in CMS. At $\sqrt{s} = 8$ TeV, the difference between ATLAS and CMS in the background-normalisation category is due to the different techniques used to estimate each background uncertainty. Additional details are discussed in appendix A.1. In the CMS tW analysis at $\sqrt{s} = 7$ TeV, the uncertainty

associated with the size of the simulated samples is evaluated as part of the total statistical uncertainty. The large difference in the pile-up uncertainty between ATLAS and CMS is due to the different methods used to assess this uncertainty, as discussed in section 5.1. At $\sqrt{s} = 8$ TeV, the sizes of the data and simulated samples used in the CMS tW analysis are smaller than in the ATLAS analysis, resulting in larger data and simulation statistical uncertainties. The large difference between the two experiments in the category of jet uncertainties arises because the JES uncertainty in ATLAS is evaluated in different categories mostly using pseudoexperiments, while in CMS the JES uncertainty is introduced as a nuisance parameter in the fit. Further details are discussed in appendix A.2. In the CMS s -channel analysis, the uncertainty associated with the size of the simulated samples is evaluated as part of the total statistical uncertainty. More details are discussed in appendix A.3.

7.1 Results

The combination of $|f_{LV}V_{tb}|^2$ is performed using the inputs from all three single-top-quark production modes. Using the same method, the combination of $|f_{LV}V_{tb}|^2$ is also performed separately for each production mode for comparison.

Combining the $|f_{LV}V_{tb}|^2$ values extracted from the t -channel and tW cross-section measurements at $\sqrt{s} = 7$ and 8 TeV from ATLAS and CMS, as well as the ATLAS s -channel measurement at $\sqrt{s} = 8$ TeV, results in

$$|f_{LV}V_{tb}|^2 = 1.05 \pm 0.02 \text{ (stat.)} \pm 0.06 \text{ (syst.)} \pm 0.01 \text{ (lumi.)} \pm 0.04 \text{ (theo.)} = 1.05 \pm 0.08,$$

with a relative uncertainty of 7.4%. The contribution from each experimental uncertainty category to the total uncertainty in the combined $|f_{LV}V_{tb}|^2$ value is shown in table 7. The theory modelling uncertainties in signal and background processes, discussed in section 5.1, dominate the experimental uncertainty and the total uncertainty. The theoretical cross-section uncertainty is the second-largest contribution to the total uncertainty in the combined $|f_{LV}V_{tb}|^2$ value. Changes in the combined $|f_{LV}V_{tb}|^2$ value from using alternative NNLO and NLO+NNLL theoretical predictions for the t -channel are less than 1%.

Figure 3 illustrates the correlations between the input measurements in the combination. The correlations are all below 0.6. The largest correlations are generally between the measurements in the same experiment at the same centre-of-mass energy, and for those that have large contributions from the same theory modelling components, such as the ATLAS s -channel measurement, which has a correlation of over 0.5 with each of the tW measurements.

The BLUE weights for each of the contributing measurements are shown in table 8. The t -channel measurements at $\sqrt{s} = 8$ TeV have the largest weight in the combination, followed by the t -channel measurements at $\sqrt{s} = 7$ TeV. The tW measurements have smaller cross-section uncertainties than the s -channel measurements, but, in addition to the correlation between tW and s -channel measurements, the tW measurements are also more correlated with the t -channel measurements in each experiment. The negative weights indicate the presence of large correlations between the corresponding measurement and some of the other measurements [81].

	<i>t</i> -channel ATLAS 8 TeV	<i>t</i> -channel CMS 8 TeV	<i>t</i> -channel ATLAS 7 TeV	<i>t</i> -channel CMS 7 TeV	<i>tW</i> ATLAS 8 TeV	<i>tW</i> CMS 8 TeV	<i>tW</i> ATLAS 7 TeV	<i>tW</i> CMS 7 TeV	<i>s</i> -channel ATLAS 8 TeV
$ f_{LV}V_{tb} ^2$	1.06	0.99	1.06	1.05	1.03	1.05	1.07	1.02	0.92
Uncertainties:									
Data statistical	0.01	0.03	0.03	0.06	0.06	0.09	0.18	0.21	0.15
Simulation statistical	0.01	0.01	0.02	0.02	0.01	0.03	0.02	–	0.11
Integrated luminosity	0.02	0.03	0.02	0.02	0.05	0.03	0.07	0.04	0.05
Theory modelling									
ISR/FSR, ren./fact. scale	0.04	0.02	0.03	0.04	0.09	0.13	0.05	0.03	0.06
NLO match., generator	0.03	0.05	0.02	0.04	0.03	–	0.11	–	0.10
Parton shower	0.02	–	–	0.01	0.02	0.15	0.16	0.10	0.02
PDF	0.01	0.02	0.03	0.01	0.01	0.02	0.02	0.02	0.03
DS/DR scheme	–	–	–	–	0.04	0.02	–	0.06	–
Top-quark p_T rew.	–	–	–	–	–	<0.01	–	–	–
Background normalisation									
Top-quark bkg.	<0.01	0.02	0.02	0.01	0.02	0.02	0.06	0.06	0.05
Other bkg. from sim.	0.01	<0.01	<0.01	0.03	0.02	0.03	0.09	0.04	0.05
Bkg. from data	<0.01	0.02	0.01	0.01	<0.01	–	0.02	–	0.01
Jets									
JES common	0.03	0.04	0.08	0.01	0.05	0.04	0.17	0.15	0.05
JES flavour	<0.01	–	0.02	–	0.02	–	–	–	0.01
JetID	<0.01	–	0.01	–	<0.01	–	0.05	–	0.01
JER	<0.01	0.01	0.02	<0.01	0.07	0.01	0.02	0.04	0.11
Detector modelling									
Leptons	0.02	0.01	0.03	0.04	0.03	0.02	0.07	0.05	0.02
HLT (had. part)	–	–	–	0.02	–	–	–	–	–
E_T^{miss} scale	<0.01	<0.01	0.03	<0.01	0.06	<0.01	–	0.03	0.01
E_T^{miss} res.	<0.01	–	–	–	<0.01	–	–	–	0.01
<i>b</i> -tagging	0.01	0.02	0.04	0.02	0.01	0.01	–	0.02	0.07
Pile-up	<0.01	0.01	<0.01	0.01	0.03	<0.01	0.11	0.01	0.01
Top-quark mass	0.01	<0.01	0.01	–	0.05	0.05	–	–	–
Theoretical cross-section									
PDF+ α_s	0.03	0.03	0.04	0.04	0.06	0.07	0.08	0.07	0.03
Ren./fact. scale	0.03	0.03	0.03	0.03	0.03	0.03	0.03	0.03	0.02
Top-quark mass	0.01	0.01	0.01	0.01	0.02	0.02	0.02	0.02	0.02
E_{beam}	<0.01	<0.01	<0.01	<0.01	<0.01	<0.01	<0.01	<0.01	<0.01
Total systematic uncertainty	0.09	0.09	0.13	0.10	0.18	0.23	0.34	0.24	0.24
Total uncertainty	0.09	0.10	0.13	0.12	0.19	0.24	0.38	0.32	0.28

Table 6. Results of the ATLAS and CMS individual $|f_{LV}V_{tb}|^2$ determinations that are the inputs to the overall $|f_{LV}V_{tb}|^2$ combination together with their experimental uncertainties. The values of $|f_{LV}V_{tb}|^2$ may slightly differ from those published for the different analyses since in this paper the theoretical cross-sections used are those marked with [†] in table 1. Experimental uncertainties contributing less than 1% are denoted by <0.01. Entries with – mean that this uncertainty was not evaluated for this analysis. Descriptions of the background categories and of the correlations of systematic uncertainties between experiments are presented in appendix A.

Combined $ f_{LV} V_{tb} ^2$	1.05	
Uncertainty category	Uncertainty	
	[%]	$\Delta f_{LV} V_{tb} ^2$
Data statistical	1.8	0.02
Simulation statistical	0.9	0.01
Integrated luminosity	1.3	0.01
Theory modelling	4.5	0.05
Background normalisation	1.3	0.01
Jets	2.6	0.03
Detector modelling	1.6	0.02
Top-quark mass	0.7	0.01
Theoretical cross-section	4.3	0.04
Total syst. unc. (excl. lumi.)	7.1	0.07
Total syst. unc. (incl. lumi.)	7.2	0.08
Total uncertainty	7.4	0.08

Table 7. Contributions from each experimental and theoretical uncertainty category to the overall $|f_{LV} V_{tb}|^2$ combination. The total uncertainty is computed by adding in quadrature all of the individual systematic uncertainties (including the integrated luminosity and theoretical cross-section) and the statistical uncertainty in data.

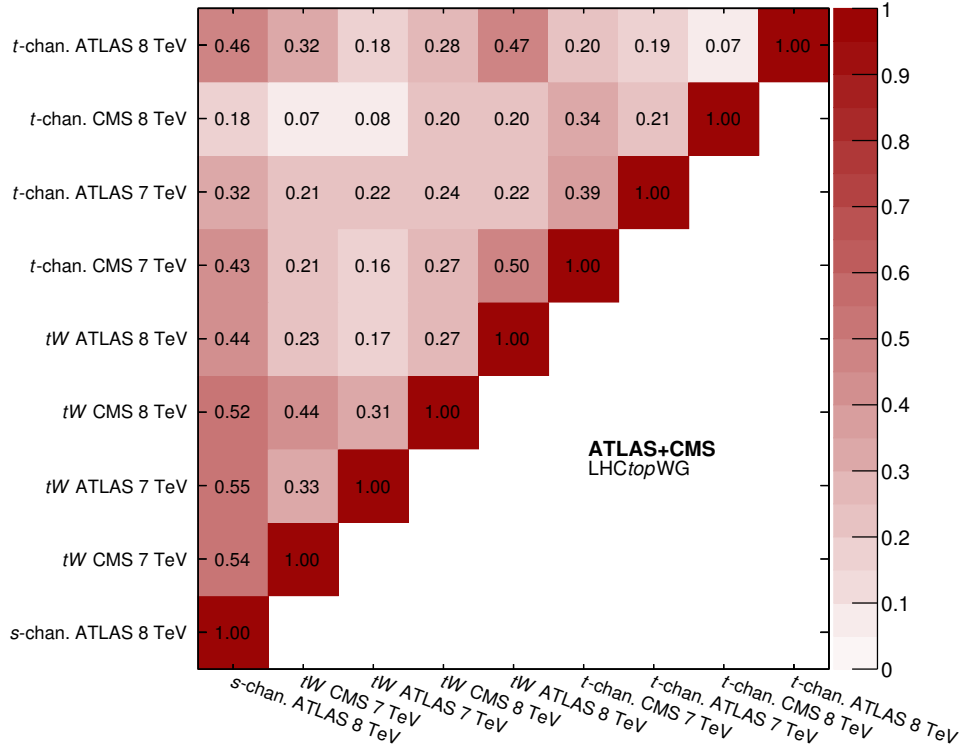


Figure 3. Correlation matrix of the overall $|f_{LV} V_{tb}|^2$ combination. Each bin corresponds to a measurement in a given production mode, experiment, and at a given centre-of-mass energy.

Process	\sqrt{s}	Experiment	BLUE weight
t -channel	8 TeV	ATLAS	0.56
		CMS	0.27
	7 TeV	ATLAS	0.07
		CMS	0.15
tW	8 TeV	ATLAS	0.05
		CMS	−0.04
	7 TeV	ATLAS	−0.02
		CMS	0.02
s -channel	8 TeV	ATLAS	−0.07

Table 8. BLUE weights for the overall $|f_{LV}V_{tb}|^2$ combination.

The combined $|f_{LV}V_{tb}|$ value from the cross-section measurements at $\sqrt{s} = 7$ and 8 TeV, including uncertainties in $\sigma_{\text{theo.}}$ for each production mode, is

$$\begin{aligned}
 |f_{LV}V_{tb}| &= 1.02 \pm 0.01 \text{ (stat.)} \pm 0.03 \text{ (syst.)} \pm 0.01 \text{ (lumi.)} \pm 0.02 \text{ (theo.)} \\
 &= 1.02 \pm 0.04 \text{ (meas.)} \pm 0.02 \text{ (theo.)} = 1.02 \pm 0.04,
 \end{aligned}$$

with a relative uncertainty of 3.7%, which improves on the precision of 4.7% of the most precise individual $|f_{LV}V_{tb}|$ extraction, which comes from the ATLAS t -channel analysis at $\sqrt{s} = 8$ TeV [32]. This is a 30% improvement over the Tevatron combination [30].

The $|f_{LV}V_{tb}|$ values are also combined for each production mode, combining across experiments and centre-of-mass energies. For the s -channel, the ATLAS and CMS measurements at $\sqrt{s} = 8$ TeV are combined. The results are

$$\begin{aligned}
 t\text{-channel} : |f_{LV}V_{tb}| &= 1.02 \pm 0.01 \text{ (stat.)} \pm 0.03 \text{ (syst.)} \pm 0.01 \text{ (lumi.)} \pm 0.02 \text{ (theo.)} \\
 &= 1.02 \pm 0.04 \text{ (meas.)} \pm 0.02 \text{ (theo.)} = 1.02 \pm 0.04, \\
 tW : |f_{LV}V_{tb}| &= 1.02 \pm 0.03 \text{ (stat.)} \pm 0.07 \text{ (syst.)} \pm 0.02 \text{ (lumi.)} \pm 0.04 \text{ (theo.)} \\
 &= 1.02 \pm 0.09 \text{ (meas.)} \pm 0.04 \text{ (theo.)} = 1.02 \pm 0.09, \\
 s\text{-channel} : |f_{LV}V_{tb}| &= 0.97 \pm 0.08 \text{ (stat.)} \pm 0.12 \text{ (syst.)} \pm 0.02 \text{ (lumi.)} \pm 0.02 \text{ (theo.)} \\
 &= 0.97 \pm 0.15 \text{ (meas.)} \pm 0.02 \text{ (theo.)} = 0.97 \pm 0.15.
 \end{aligned}$$

The relative uncertainties are 3.9%, 8.4% and 15.0% respectively. In all cases, these results are more precise than the best individual determinations of $|f_{LV}V_{tb}|$, which have uncertainties of 4.7%, 9.9% and 20.8% for the t -channel [32], tW [34] and s -channel [37] analyses respectively.

Figure 4 shows a summary of the $|f_{LV}V_{tb}|$ combinations. The combination is dominated by the t -channel analyses.

7.2 Stability tests

The stability of the combination of the $|f_{LV}V_{tb}|^2$ values to variations in the correlation assumptions, discussed in section 5, is checked for the dominant uncertainty contributions.

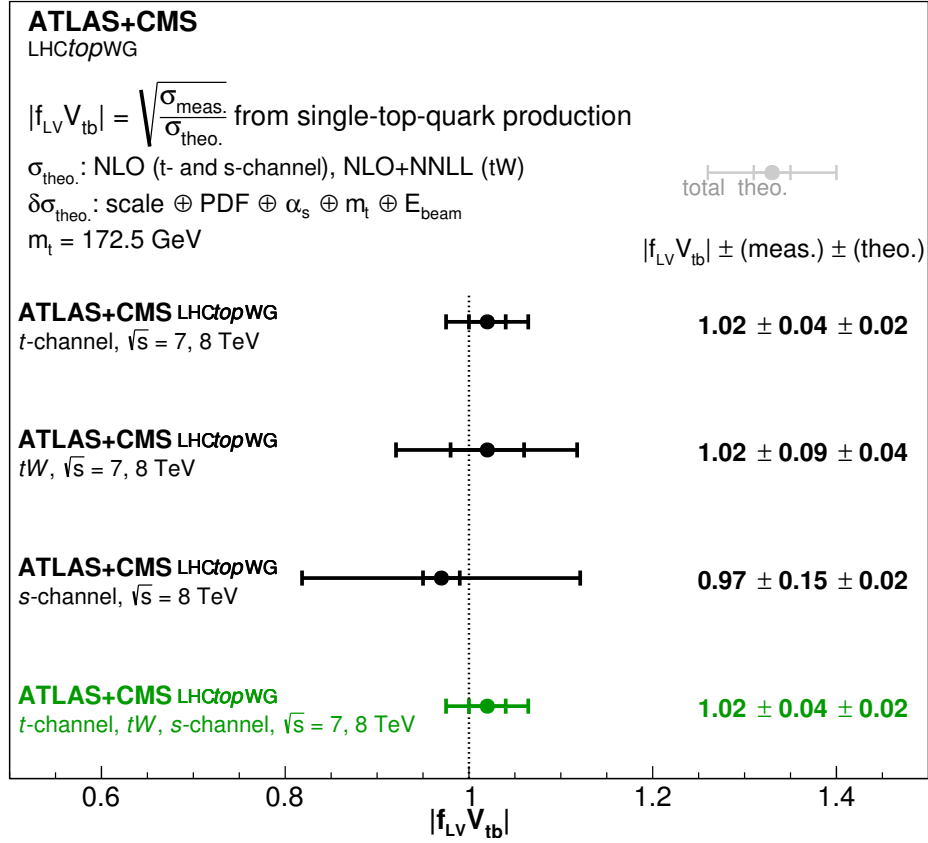


Figure 4. The combined $|f_{LV}V_{tb}|$ value extracted from the t -channel and tW cross-section measurements at $\sqrt{s} = 7$ and 8 TeV from ATLAS and CMS, as well as the ATLAS s -channel measurement at $\sqrt{s} = 8$ TeV, is shown together with the combined $|f_{LV}V_{tb}|$ values for each production mode. The theoretical predictions for t -channel and s -channel production are computed at NLO accuracy, while the theoretical predictions for tW are calculated at NLO+NNLL accuracy. The $\sigma_{\text{theo.}}$ uncertainties used to compute $|f_{LV}V_{tb}|$ include scale, PDF+ α_s , m_t , and E_{beam} variations.

The correlation values are varied for the theory modelling, JES, and the most important contributions to the theoretical cross-section predictions (i.e. PDF+ α_s and scale). Because of the scheme that is used for the correlations, stability tests are also performed for the uncertainties associated with the integrated luminosity. Figure 5 summarises the results of these stability tests, where the correlations between ATLAS and CMS (and also between centre-of-mass energies for the integrated luminosity) are varied.

The uncertainties in the theory modelling category (i.e. scales and radiation modelling, NLO matching, and PS and hadronisation) are varied from their default value of fully correlated to half correlated and to the more extreme tests of uncorrelated and half anti-correlated. The JES category is varied from its default value of uncorrelated to half correlated and half anti-correlated and the more extreme variation of fully correlated. The theoretical cross-section uncertainties, PDF+ α_s and scale, are varied from their default values of fully correlated to half correlated, uncorrelated and half anti-correlated. For the integrated luminosity, the correlation between ATLAS and CMS is varied from its default

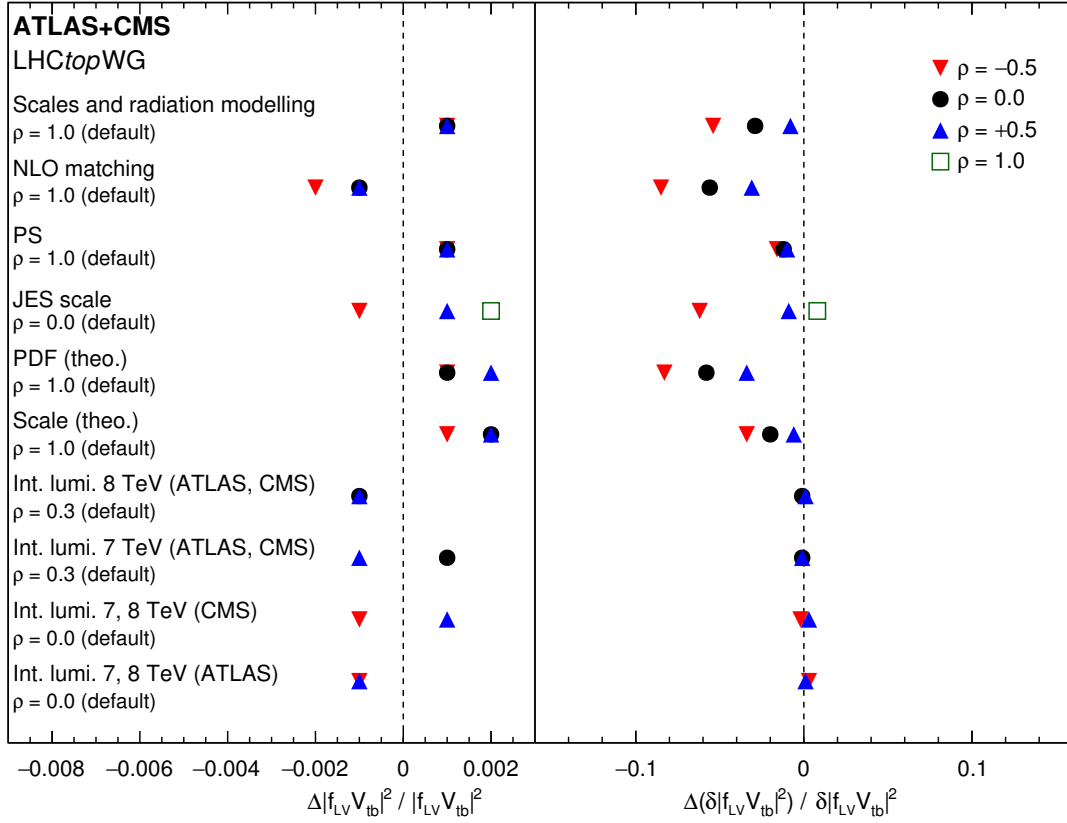


Figure 5. Results of the stability tests performed by varying of the correlation assumptions in different uncertainty categories: theory modelling (scales and radiation modelling, NLO matching, and PS and hadronisation), JES, dominant theoretical cross-section predictions (i.e. PDF+ α_s and scale) and integrated luminosity. Two or three variations are considered depending on the uncertainty category. The corresponding relative shifts (with shift = varied – nominal) in the central value, $\Delta|f_{LV}V_{tb}|^2/|f_{LV}V_{tb}|^2$, and in its uncertainty, $\Delta(\delta|f_{LV}V_{tb}|^2)/(\delta|f_{LV}V_{tb}|^2)$, are shown.

value of 30% correlated to half correlated and uncorrelated. The correlation between different centre-of-mass energies for each experiment is varied from the default of uncorrelated to half and fully correlated. The correlation of the theoretical scale uncertainty between different processes is also tested. For all variations, the relative changes in the central value of the combined $|f_{LV}V_{tb}|$ are significantly smaller ($<0.5\%$) than the relative total uncertainty of 3.7%. Additionally, the relative changes in the total uncertainty are below 0.004, i.e., less than 10% of the total uncertainty of 0.04. These tests show that the result of the combination is robust and does not critically depend on any of the correlation assumptions. The cross-section combinations similarly do not depend significantly on any of the correlation assumptions.

8 Summary

The combinations of single-top-quark production cross-section measurements in the t -channel, tW , and s -channel production modes are presented, using data from LHC pp

collisions collected by the ATLAS and CMS Collaborations. The combinations for each production mode are performed at $\sqrt{s} = 7$ and 8 TeV, using data corresponding to integrated luminosities of 1.17 to 5.1 fb $^{-1}$ at $\sqrt{s} = 7$ TeV, and of 12.2 to 20.3 fb $^{-1}$ at $\sqrt{s} = 8$ TeV. The combined t -channel cross-sections are found to be 67.5 ± 5.7 pb and 87.7 ± 5.8 pb at $\sqrt{s} = 7$ and 8 TeV respectively. The values of the combined tW cross-sections at $\sqrt{s} = 7$ and 8 TeV are 16.3 ± 4.1 pb and 23.1 ± 3.6 pb respectively. For the s -channel cross-section, the combination yields 4.9 ± 1.4 pb at $\sqrt{s} = 8$ TeV. The square of the magnitude of the CKM matrix element V_{tb} multiplied by a form factor accounting for possible contributions from physics beyond the SM, f_{LV} , is determined from each production mode at each centre-of-mass energy, using the ratio of the measured cross-section to its theoretical prediction, and assuming that the top-quark-related CKM matrix elements obey the relation $|V_{td}|, |V_{ts}| \ll |V_{tb}|$. The values of $|f_{LV}V_{tb}|^2$ extracted from individual ratios at $\sqrt{s} = 7$ and 8 TeV yield a combined value of $|f_{LV}V_{tb}| = 1.02 \pm 0.04$ (meas.) ± 0.02 (theo.). All combined measurements are consistent with their corresponding SM predictions.

A Systematic uncertainties in cross-section measurements

The single-top-quark cross-sections measured by the ATLAS and CMS Collaborations at $\sqrt{s} = 7$ and 8 TeV, as well as the uncertainties and their correlations between experiments, are summarised in tables 9–13 for the t -channel, tW , and s -channel production modes. Similar to the approach that is followed in combinations using the BLUE method, the total uncertainty in these tables is evaluated as the sum in quadrature of the individual uncertainties. To obtain the impact of each source of uncertainty, the input analyses use either pseudoexperiments or approximate procedures which neglect the correlations between sources of uncertainty introduced by the fit to data. In the latter case, this may lead to small changes in the total uncertainty compared with the input measurements presented in table 2. The likelihood fit includes all nuisance parameters at the same time to evaluate the total uncertainty. The method used by each input analysis to evaluate the individual uncertainties is described below.

A.1 Systematic uncertainties in t -channel cross-section measurements

The t -channel cross-sections measured by the ATLAS and CMS Collaborations at $\sqrt{s} = 7$ TeV [59, 60] and $\sqrt{s} = 8$ TeV [32, 33], as well as the uncertainties and their correlations between experiments, are shown in tables 9 and 10 respectively. The total uncertainty given for each measurement is the sum in quadrature of the individual uncertainties. This is slightly different from the total uncertainty shown in table 2 for the CMS measurements at $\sqrt{s} = 7$ and 8 TeV since the total uncertainty is evaluated, through the fit, by varying all nuisance parameters at the same time.

In table 9, the CMS result at $\sqrt{s} = 7$ TeV has a larger data statistical uncertainty than the ATLAS result because the two experiments use data samples of different size (see table 2). In the background-normalisation category, the “Bkg. from MC” uncertainty refers to the tW , s -channel, $t\bar{t}$, W/Z +jets, and diboson backgrounds. In the ATLAS measurement, the normalisation uncertainty in the multijet background is estimated by comparing

	ATLAS ($\sigma_{t\text{-chan.}}, \sqrt{s} = 7\text{ TeV}$)		CMS ($\sigma_{t\text{-chan.}}, \sqrt{s} = 7\text{ TeV}$)		
Cross-section	68.0 pb		67.2 pb		
Uncertainty category	Uncertainty		Uncertainty		ρ
Data statistical		2.7%		5.8%	0.0
Simulation statistical		1.9%		1.9%	0.0
Integrated luminosity		1.8%		2.2%	0.3
Theory modelling	Ren./fact. scales, ISR/FSR	2.6%	Ren./fact. scales	3.5%	1.0
	NLO match., PS ($t\bar{t}$, t -chan.)	2.2%	Sig. modelling (NLO method)	4.3%	1.0
			Parton shower	0.8%	1.0
	PDF	3.2%	PDF	1.4%	1.0
Category subtotal	4.7%		5.8%		0.85
Background norm.	Bkg. from MC: norm.	1.6%	Bkg. from MC: norm.	2.7%	0.0
	Bkg. from MC/data: multijet norm.	1.4%	Bkg. from data: multijet norm.	1.3%	0.0
Category subtotal	2.1%		3.0%		0.0
Jets	JES common	7.6%	JES	0.9%	0.0
	JES flavour	1.8%			0.0
	JetID	1.1%			0.0
	JER	1.9%	JER	0.3%	0.0
Category subtotal	8.1%		0.9%		0.0
Detector modelling	Lepton modelling	2.8%	Lepton modelling	3.5%	0.0
			HLT (had. part)	1.5%	0.0
	E_T^{miss} modelling	2.6%	E_T^{miss} modelling	0.1%	0.0
	b -tagging	3.9%	b -tagging	2.2%	0.0
	Pile-up	0.2%	Pile-up	0.6%	0.0
Category subtotal	5.5%		4.4%		0.0
Total uncertainty	11.7%		10.2%		0.20

Table 9. Measured cross-sections, uncertainty components, their magnitudes (relative to the individual measurements) and the correlation (ρ) between the ATLAS and CMS $\sigma_{t\text{-chan.}}$ measurements at $\sqrt{s} = 7\text{ TeV}$. Uncertainties in the same row can be compared between experiments, as detailed in the text. The naming conventions follow those of the corresponding experiments.

background estimates made using different techniques based on data and simulation samples, while in the CMS measurement, it is estimated from the difference between alternative methods based on data. There is also a large difference between the two experiments in the jets category. As discussed in section 5.1, the uncertainty in each JES component in the ATLAS measurement is evaluated using pseudoexperiments. The CMS measurement is a BLUE combination of three different measurements, two of which introduce JES components as a nuisance parameter in the fit. Since these fits use additional control regions, the impact of the JES is reduced. In addition, the JES uncertainty in the analyses at $\sqrt{s} = 7\text{ TeV}$ is smaller for CMS [100] than for ATLAS [99].

In the analyses at $\sqrt{s} = 8\text{ TeV}$, summarised in table 10, the difference between ATLAS and CMS in the background normalisation category is due to the different techniques used to estimate each background uncertainty. The “Other bkg. from MC” uncertainty includes the contributions from the tW , s -channel, $t\bar{t}$, W/Z +jets, and diboson backgrounds in the ATLAS analysis, and the tW , s -channel, Z +jets, and diboson backgrounds in the CMS analysis. In the ATLAS measurement, the normalisation uncertainties associated with the

	ATLAS ($\sigma_{t\text{-chan.}}, \sqrt{s} = 8 \text{ TeV}$)		CMS ($\sigma_{t\text{-chan.}}, \sqrt{s} = 8 \text{ TeV}$)		
Cross-section	89.6 pb		83.6 pb		
Uncertainty category	Uncertainty		Uncertainty		ρ
Data statistical		1.4%		2.7%	0.0
Simulation statistical		0.8%		0.7%	0.0
Integrated luminosity		1.9%		2.6%	0.3
Theory modelling	Ren./fact. scales	3.6%	Ren./fact. scales	1.9%	1.0
	NLO match.	3.3%	NLO match., 4FS vs 5FS	4.9%	1.0
	Parton shower	2.1%			1.0
	PDF	1.3%	PDF	1.9%	1.0
Category subtotal		5.5%		5.6%	0.84
Background norm.	$t\bar{t}$, tW and s -chan. norm.	0.1%	$t\bar{t}$ and W +jets norm.	2.2%	0.0
	Other bkg. from MC: norm.	0.9%	Other bkg. from MC: norm.	0.3%	0.0
	Bkg. from MC/data: multijet norm.	0.3%	Bkg. from data: multijet norm.	2.3%	0.0
Category subtotal		1.0%		3.2%	0.0
Jets	JES common	3.2%	JES	4.2%	0.0
	JES flavour	0.2%			0.0
	JetID	0.1%			0.0
	JER	0.4%	JER	0.7%	0.0
Category subtotal		3.2%		4.3%	0.0
Detector modelling	Lepton modelling	1.9%	Lepton modelling	0.6%	0.0
	E_T^{miss} scale	0.4%	E_T^{miss} modelling	0.3%	0.0
	E_T^{miss} resolution	0.2%			0.0
	b -tagging	1.1%	b -tagging	2.5%	0.0
	Pile-up	0.3%	Pile-up	0.7%	0.0
Category subtotal		2.3%		2.7%	0.0
Total uncertainty		7.3%		9.0%	0.42

Table 10. Measured cross-sections, uncertainty components, their magnitudes (relative to the individual measurements) and the correlation (ρ) between the ATLAS and CMS $\sigma_{t\text{-chan.}}$ measurements at $\sqrt{s} = 8 \text{ TeV}$. Uncertainties in the same row can be compared between experiments, as detailed in the text. The naming conventions follow those of the corresponding experiments.

top-quark, W/Z +jets, diboson, and multijet backgrounds are estimated using pseudoexperiments. Variations in the theoretical cross-section predictions for these processes are also considered, except for the multijet background, where the results obtained from data and simulation samples analysed with various techniques are compared. In the CMS measurement, the uncertainty in the multijet background is estimated from the difference between alternative methods based on data. The normalisations of the $t\bar{t}$ and W +jets backgrounds are included as nuisance parameters in the fit, while the shapes of their distributions are adjusted by corrections based on data in control regions.

A.2 Systematic uncertainties in tW cross-section measurements

The tW cross-sections measured by the ATLAS and CMS Collaborations at $\sqrt{s} = 7 \text{ TeV}$ [61, 62] and $\sqrt{s} = 8 \text{ TeV}$ [34, 35], as well as the uncertainties and their correlations between experiments, are shown in tables 11 and 12 respectively.

In table 11, the CMS measurement at $\sqrt{s} = 7 \text{ TeV}$ takes into account the uncertainty associated with the size of the simulated event samples using the Barlow-Beeston

	ATLAS (σ_{tW} , $\sqrt{s} = 7$ TeV)		CMS (σ_{tW} , $\sqrt{s} = 7$ TeV)		
Cross-section	16.8 pb		16.0 pb		
Uncertainty category	Uncertainty		Uncertainty		ρ
Data statistical		17.0%		20.8%	0.0
Simulation statistical		2.0%		0.0%	0.0
Integrated luminosity		7.0%		4.3%	0.3
Theory modelling	ISR/FSR, scales	5.0%	ISR/FSR, scales	2.8%	1.0
	$tW/t\bar{t}$ NLO match.	10.0%			1.0
	$tW/t\bar{t}$ PS	15.0%	tW ME/PS match. thr.	10.1%	1.0
	PDF	2.0%	PDF	2.1%	1.0
			DR/DS scheme	5.9%	1.0
Category subtotal	18.8%		12.2%		0.74
Background norm.	$t\bar{t}$ norm.	6.0%	$t\bar{t}$ norm.	6.0%	0.0
	Z +jets, diboson norm.	8.0%	Z/γ^* +jets norm.	4.2%	0.0
	Bkg. from data: fake lept. norm.	2.0%			0.0
Category subtotal	10.2%		7.3%		0.0
Jets	JES	16.0%	JES	15.1%	0.0
	JetID	5.0%			0.0
	JER	2.0%	JER	3.6%	0.0
Category subtotal	16.9%		15.6%		0.0
Detector modelling	Lepton modelling	7.0%	Lepton modelling	5.2%	0.0
			E_T^{miss} modelling	2.5%	0.0
			b -tagging	1.9%	0.0
	Pile-up	10.0%	Pile-up	1.5%	0.0
Category subtotal	12.2%		6.2%		0.0
Total uncertainty	35.1%		30.6%		0.17

Table 11. Measured cross-sections, uncertainty components, their magnitudes (relative to the individual measurements) and the correlation (ρ) between the ATLAS and CMS σ_{tW} measurements at $\sqrt{s} = 7$ TeV. Uncertainties in the same row can be compared between experiments, as detailed in the text. The naming conventions follow those of the corresponding experiments.

method [125]. This contribution is included as part of the total statistical uncertainty. This uncertainty is therefore considered to be zero for the CMS measurement to avoid double-counting. Since the statistical uncertainties in the data and simulation are uncorrelated between the two experiments, this choice has almost no effect on the combination. In the ATLAS analysis, the normalisation uncertainty in the misidentified lepton (fake lept.) background is conservatively taken to be 100%, based on comparisons in data. The E_T^{miss} uncertainties are included in the pile-up modelling uncertainty. The b -tagging uncertainty is not considered because no b -tagging criterion is required in the event selection. The large difference in the pile-up uncertainty between ATLAS and CMS arises from different methods employed by the experiments to assess this uncertainty, as discussed in section 5.1.

In table 12, the tW measurement by CMS at $\sqrt{s} = 8$ TeV is based on the first half of the $\sqrt{s} = 8$ TeV data sample. This leads to a larger data statistical uncertainty for CMS than for ATLAS. For the same reason, the sizes of the simulated event samples are smaller, resulting in a larger simulation statistical uncertainty in the CMS result. In the ATLAS measurement, the normalisation uncertainty in the multijet background is estimated by comparing estimates made using different techniques on data and simulation

	ATLAS (σ_{tW} , $\sqrt{s} = 8$ TeV)		CMS (σ_{tW} , $\sqrt{s} = 8$ TeV)		
Cross-section	23.0 pb		23.4 pb		
Uncertainty category	Uncertainty		Uncertainty		ρ
Data statistical		5.8%		8.1%	0.0
Simulation statistical		0.5%		2.4%	0.0
Integrated luminosity		4.6%		3.0%	0.3
Theory modelling	ISR/FSR	8.8%	Ren./fact. scales	12.4%	1.0
	NLO match.	2.5%			1.0
	Parton shower	1.7%	Parton shower	14.1%	1.0
	PDF	0.6%	PDF	1.7%	1.0
	$tW/t\bar{t}$ overlap	3.5%	tW DR/DS scheme	2.1%	1.0
			Top-quark p_T reweight.	0.4%	0.0
Category subtotal	10.0%		19.0%		0.75
Background norm.	$t\bar{t}$ norm.	1.9%	$t\bar{t}$ norm.	1.7%	0.0
	Z +jets, diboson norm.	2.0%	Z +jets norm.	2.6%	0.0
	Bkg. from data: fake lept. norm.	0.3%			0.0
Category subtotal	2.8%		3.1%		0.0
Jets	JES common	5.3%	JES	3.8%	0.0
	JES flavour	1.9%			0.0
	JetID	0.2%			0.0
	JER	6.5%	JER	0.9%	0.0
Category subtotal	8.6%		3.9%		0.0
Detector modelling	Lepton modelling	3.0%	Lepton modelling	1.8%	0.0
	E_T^{miss} scale	5.5%	E_T^{miss} modelling	0.4%	0.0
	E_T^{miss} resolution	0.2%			0.0
	b -tagging	1.0%	b -tagging	0.9%	0.0
	Pile-up	2.7%	Pile-up	0.4%	0.0
Category subtotal	6.9%		2.0%		0.0
Total uncertainty	16.8%		21.7%		0.40

Table 12. Measured cross-sections, uncertainty components, their magnitudes (relative to the individual measurements) and the correlation (ρ) between the ATLAS and CMS σ_{tW} measurements at $\sqrt{s} = 8$ TeV. Uncertainties in the same row can be compared between experiments, as detailed in the text. The naming conventions follow those of the corresponding experiments.

samples, while in the CMS measurement, the uncertainty contribution of the multijet background is estimated from the difference between alternative methods based on data. In the ATLAS analysis, the misidentified lepton and non-prompt (fake lept.) background has a normalisation uncertainty of 60%, based on comparisons in data, to account for possible mismodelling of the jet multiplicity and jet acceptance. There is a large difference between the two experiments in the jets category. As discussed in section 5.1, the JES uncertainty in the ATLAS measurement is evaluated in different categories. The detector modelling component of the JES common uncertainty is constrained in the fit to data. In the CMS measurement, different components of the JES uncertainty are grouped together, and the group is introduced as a nuisance parameter in the fit. The E_T^{miss} modelling uncertainty is smaller for the CMS measurement due to the use of low- p_T jets, which allows this uncertainty to be constrained in the fit to data, as discussed in section 5.1.

	ATLAS ($\sigma_{s\text{-chan.}}, \sqrt{s} = 8 \text{ TeV}$)		CMS ($\sigma_{s\text{-chan.}}, \sqrt{s} = 8 \text{ TeV}$)		
Cross-section	4.8 pb		13.4 pb		
Uncertainty category	Uncertainty		Uncertainty		ρ
Data statistical		16.0%		10.0%	0.0
Simulation statistical		12.0%		0.0%	0.0
Integrated luminosity		5.0%		4.0%	0.3
Theory modelling	Ren./fact. scales	7.0%	Ren./fact. scales	30.0%	1.0
	$t\bar{t}$, t -chan. generator	11.0%			1.0
	Parton shower	2.0%	Parton shower	7.0%	1.0
	PDF	3.0%	PDF	7.0%	1.0
			Top-quark p_T reweight.	6.0%	0.0
Category subtotal	13.5%		32.2%		0.56
Background norm.	t -chan., $t\bar{t}$ norm.	5.0%	t -chan., $t\bar{t}$ norm.	12.0%	0.0
	W/Z +jets, diboson norm.	6.0%	W/Z +jets, diboson norm.	12.0%	0.0
	Bkg. from data: multijet norm.	1.0%	Bkg. from data: multijet norm.	2.0%	0.0
Category subtotal	7.9%		17.1%		0.0
Jets	JES common	5.0%	JES	32.5%	0.0
	JES flavour	1.0%			0.0
	JetID	1.0%			0.0
	JER	12.0%	JER	10.2%	0.0
Category subtotal	13.1%		34.1%		0.0
Detector modelling	Lepton modelling	2.4%	Lepton modelling	1.0%	0.0
	E_T^{miss} scale	1.0%	E_T^{miss} modelling	6.0%	0.0
	E_T^{miss} res	1.0%			0.0
	b -tagging	8.0%	b -tagging	14.0%	0.0
	Pile-up	1.0%	Pile-up	9.0%	0.0
Category subtotal	8.5%		17.7%		0.0
Total uncertainty	30.2%		54.0%		0.15

Table 13. Measured cross-sections, uncertainty components, their magnitudes (relative to the individual measurements) and the correlation (ρ) between the ATLAS and CMS $\sigma_{s\text{-chan.}}$ measurements at $\sqrt{s} = 8 \text{ TeV}$. Uncertainties in the same row can be compared between experiments, as detailed in the text. The naming conventions follow those of the corresponding experiments.

A.3 Systematic uncertainties in s -channel cross-section measurements

The s -channel cross-sections measured by the ATLAS and CMS Collaborations at $\sqrt{s} = 8 \text{ TeV}$ [37, 63], as well as the uncertainties and their correlations between experiments, are shown in table 13.

The CMS measurement takes into account the uncertainty associated with the size of the simulated event samples using the Barlow-Beeston method [125]. The contribution is included in the total statistical uncertainty. This uncertainty is therefore considered to be zero for the CMS measurement to avoid double-counting. Since the statistical uncertainties in the data and simulation are uncorrelated between the two experiments, this choice has almost no effect on the combination. The result from ATLAS has smaller uncertainties. This is attributed to the use of the latest simulation samples with tuned parameters [126] as well as the use of the matrix element method in the ATLAS analysis. In addition, all systematic uncertainties are profiled in the ATLAS analysis, while in the CMS analysis, major

uncertainties, including those from jets and in the theory modelling category, are excluded from the fit and evaluated using pseudoexperiments. The total uncertainties in table 13 are slightly different from the uncertainties shown in table 2 because here the uncertainties are summed in quadrature, while in the input analyses the impacts of at least some of the uncertainties are included in the fits to data. In particular, the difference between the relative total uncertainty shown in table 2 for the ATLAS measurement, i.e. 34.4%, and the relative total uncertainty shown in table 13 is due to the usage of an approximate procedure to compute the individual uncertainty contributions. Possible correlation terms between the systematic uncertainties introduced by the fit are not included here.

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The ATLAS collaboration

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